

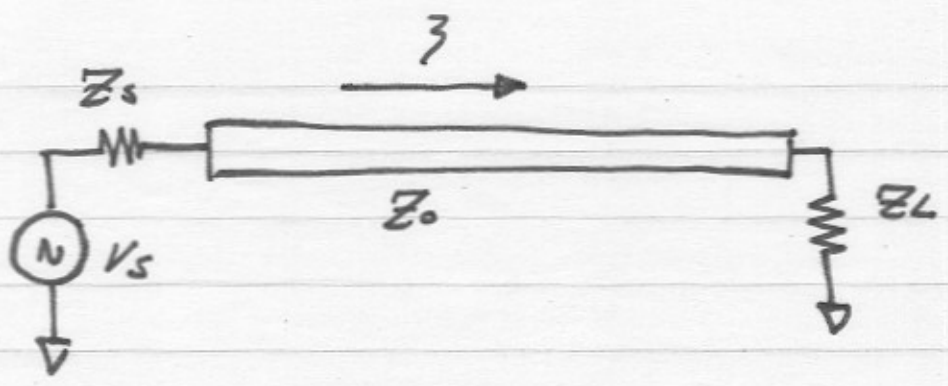
ECE202A Notes set 2

Transmission Lines in the frequency domain
and the smith chart.

Time-domain analysis is intellectually clearer,
the picture being forward & reverse waves
propagating, reflecting, and re-reflecting.

This analysis becomes intractable as soon as
we introduce reactive ($-j\omega L$ & $-j/\omega C$) impedances
as multiple convolutions will be required for
time-domain reflection analysis.

→ Analyze in the frequency domain instead
this is the classical approach.



assume: $V_s = \text{Re} \{ V_0 e^{-j\omega t} \}$

V_0 is complex: $V_0 = \|V_0\| \cdot \exp\{-j\phi_0\}$

so $v_s(t) = \|V_0\| \cdot \cos(\omega t + \phi_0)$

on transmission line, waves travel as

$(z \pm vt)$ or equivalently $(t \pm z/v)$

For a wave traveling at velocity v ,

$\cos(\omega t + \phi_0) \rightarrow \cos(\omega(t \pm z/v) + \phi_0)$

$= \cos(\omega t + \phi_0 \pm z \cdot \omega/v)$

$= \cos(\omega t + \phi_0 \pm \beta z)$

where β is the phase constant $\beta = \omega/v$

$$\beta = \omega/v = 2\pi/\lambda$$

... where λ is the wavelength

For exponential waves, the form will thus be

$$v_0 e^{-j\omega t} e^{-j\phi_0} e^{\pm j\beta z}$$

and, as always with phasor notation, the $e^{-j\omega t}$ time dependence is taken as implicit:

Voltage on the transmission line:

$$V(z) = V^+(z) + V^-(z)$$

$$= V^+(0) e^{-j\beta z} + V^-(0) e^{+j\beta z}$$

$$Z_0 I(z) = V^+(z) - V^-(z)$$

$$= V^+(0) e^{-j\beta z} - V^-(0) e^{+j\beta z}$$

Wave parameters, again:

define wave amplitudes such that if $a=1$
then wave power = 1 watt

$\rightarrow a(z) = V_0^+(z) / \sqrt{Z_0}$ Forward wave amplitude

$b(z) = V^-(z) / \sqrt{Z_0}$ Reverse wave amplitude.

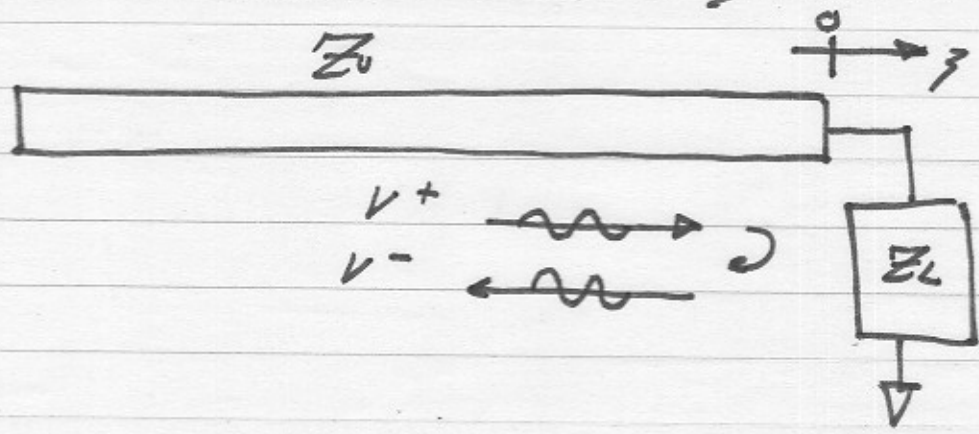
Power in Forward wave = aa^*

" " reverse " = bb^*

It is assumed throughout that we are using R.M.S. quantities.

* Note that the relationships must be generalized if Z_0 is complex.

Reflections in the Frequency domain:

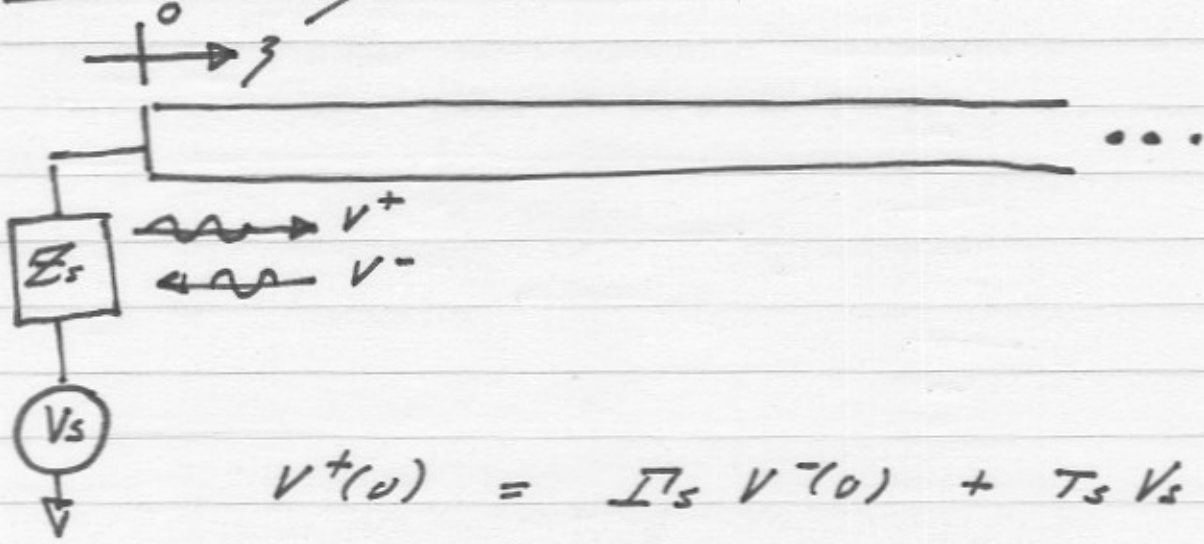


$$V^-(z) = V^+(z) \Gamma_L$$

$$\Gamma_L = \frac{Z_L - Z_0}{Z_L + Z_0} \quad \text{where}$$

$Z_L = Z_L / Z_0$ is the normalized load impedance

From a generator:



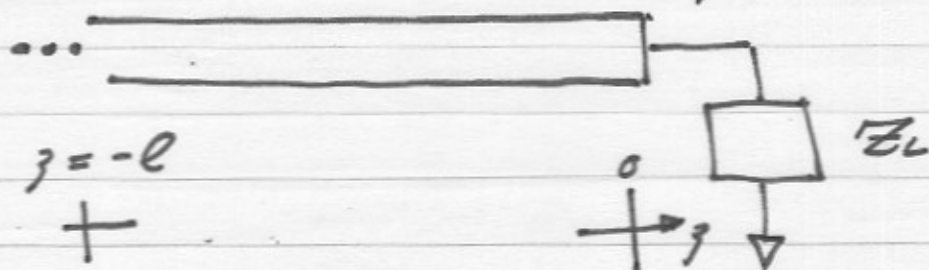
T_s = source transmission coefficient

$$= \frac{Z_0}{Z_s + Z_0}$$

$$\Gamma_s = \frac{\zeta_s - 1}{\zeta_s + 1} \quad \text{source reflection coefficient}$$

$$\zeta_s = Z_s / Z_0 \quad \text{normalized source impedance.}$$

Movement of reference plane



$$V(z) = V^+(z) + V^-(z)$$

$$= V^+(z) [1 + \Gamma(z)]$$

where $\Gamma(z) \triangleq V^-(z)/V^+(z)$

is the position-dependent impedance reflection-coefficient

$$V(z) = V^+(0) e^{-j\beta z} [1 + \Gamma(0) e^{+2j\beta z}]$$

because

$$\Gamma(z) = \frac{V^-(z)}{V^+(z)} = \frac{V^-(0) e^{+j\beta z}}{V^+(0) e^{-j\beta z}} = \Gamma(0) e^{2j\beta z}$$

So:

$\Gamma(-l)$ = the reflection coefficient at a distance l from the load

$$= \Gamma(0) e^{-2j\beta l}$$

So, the reflection coefficient has gone through

a phase shift of

$$\underline{\text{minus}} \quad \frac{l}{\lambda} \cdot 2 \cdot 2\pi \quad \text{radians}$$

or

$$\underline{\text{minus}} \quad 2\beta l \quad \text{radians}$$

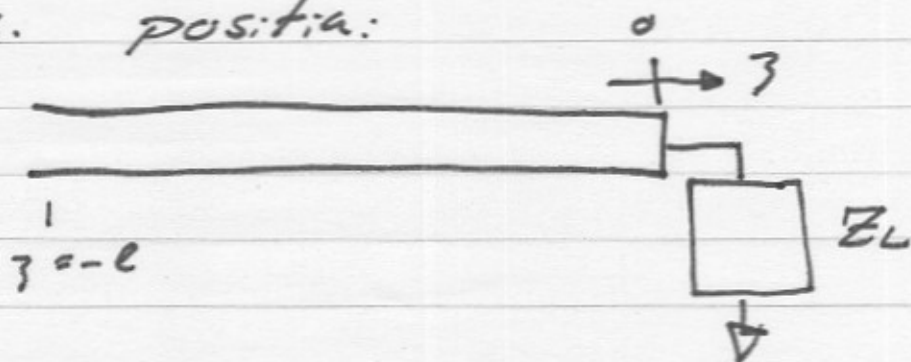
or

$$\underline{\text{minus}} \quad \frac{l}{\lambda} \cdot 2 \cdot 360 \quad \text{degrees.}$$

... reflection coefficient changes phase with position.

Impedance vs. position:

same picture:



Impedance at ~~z=l~~ any point:

$$Z(z) \triangleq V(z) / I(z)$$

$$= [V^+(z) + V^-(z)] / \left[\frac{1}{Z_0} (V^+(z) - V^-(z)) \right]$$

$$= Z_0 \frac{1 + \Gamma(z)}{1 - \Gamma(z)}$$

Normalized impedance at any point

$$\tilde{z}(z) \triangleq Z(z) / Z_0 = \frac{1 + \Gamma(z)}{1 - \Gamma(z)}$$

Input impedance of line at $z = -l$

$$\tilde{z}(-l) = \frac{1 + \Gamma(-l)}{1 - \Gamma(-l)} \quad \text{normalized}$$

$$Z(-l) = Z_0 \tilde{z}(-l) \quad \text{unnormalized.}$$

So: Impedance depends upon position.

Summarizing:

$$\left. \begin{aligned} V(z) &= V^+(z) + V^-(z) \\ Z_0 I(z) &= V^+(z) - V^-(z) \end{aligned} \right\} \text{waves}$$

$$\left. \begin{aligned} \Gamma(z) &= V^-(z) / V^+(z) \\ &= \Gamma(0) e^{2j\beta z} \end{aligned} \right\} \text{reflection coefficient}$$

$$\zeta(z) = \frac{1 + \Gamma(z)}{1 - \Gamma(z)} \quad \left. \vphantom{\zeta(z)} \right\} \text{normalized impedance.}$$

while this is conceptually very simple, it is a great deal of math.

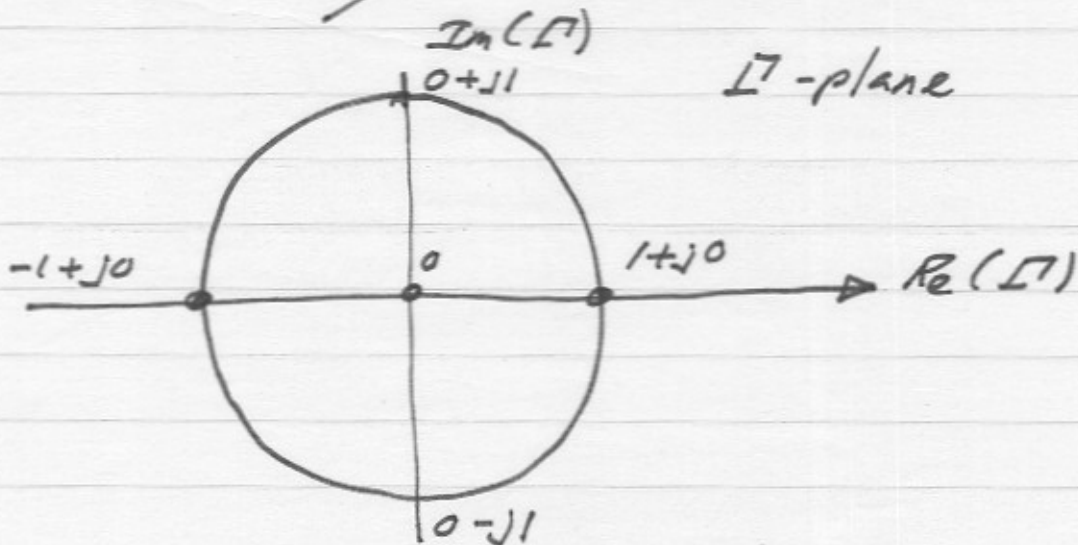
⇒ A graphical representation would help
this is the Smith chart.

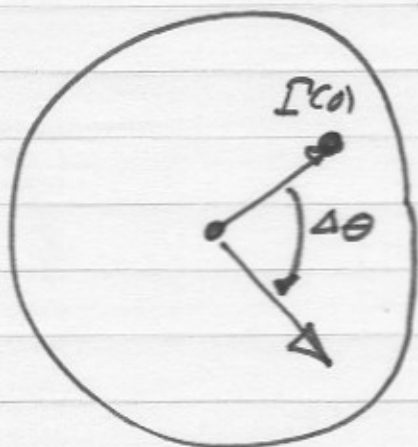
The relationship

$$\Gamma = \frac{1 + \Gamma}{1 - \Gamma} \quad \Gamma \text{ and } \Gamma \text{ both complex.}$$

is key. It is a 1:1 mapping between complex numbers, and is in fact an analytic function and a conformal transformation. Review your complex analysis if necessary and interested.

In the 2-dimensional plane of values of Γ - the Γ -plane - a reflection coefficient Γ is represented by a point:



Γ -plane

as we move away from the load by a distance l on the transmission line, Γ rotates by an angle $\Delta\theta$

$$\Delta\theta = -2\beta l$$

$$= -360^\circ \times \frac{l}{\lambda} \cdot 2$$

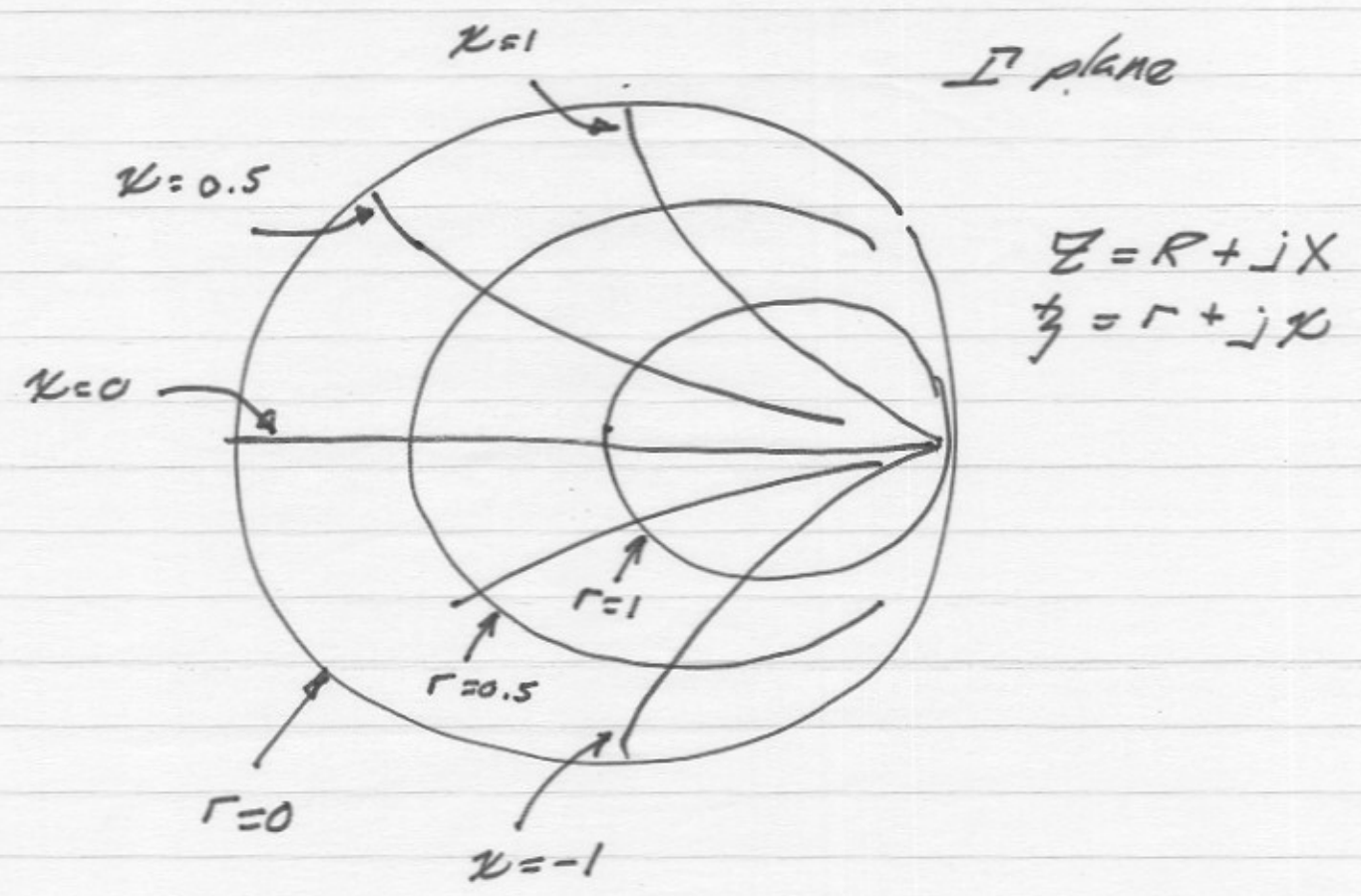
= one whole rotation in the Γ plane

For each half-wavelength movement on the line.

$$\zeta = \frac{1 + \Gamma}{1 - \Gamma}$$

This is a 1:1 relationship between Γ (magnitude & phase) and ζ (real & imaginary)

Plot the units of normalized impedance on the Γ -plane!



This is the Smith chart.

So, given some load impedance Z_L , we can calculate Γ just by looking up the Z -coordinates of the point and reading off the magnitude and phase of Γ with a straightedge and protractor.

Change of impedance vs position: just rotate the Γ -vector clockwise around the chart for the distance l , at a rate of one ^{rotation} wavelength for every half-wavelength of movement.

- Then read off impedance from the chart.

Impedance - Admittance Chart

Impedance

$$Z = R + jX$$

resistance reactance

Normalized Impedance

$$\bar{z} = Z/Z_0 = r + jx$$

Admittance

$$Y = 1/Z = G + jB$$

Conductance susceptance

Normalized admittance

$$\bar{y} = Y/Y_0 = Y \cdot Z_0 = g + jb$$

We can plot the values of

r and x

- and -

g and b

on the Γ -plane. This is the impedance - admittance chart.

Among several advantages,

$$\text{if } Z = 50 + j50 \Omega$$

$$Z_0 = 50 \Omega$$



$$\text{so that } \bar{z} = 1 + j1$$

then $Y = ?$

while we could use complex math, we can just use the plot:

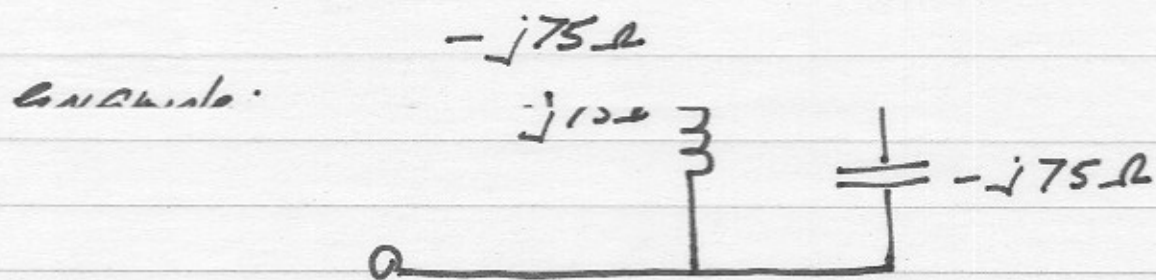
draw the point $\bar{z} = 1 + j1$ and read off the normalized admittance as $y = 0.5 - j0.5$

un-normalize: $Y = \frac{0.5 - j0.5}{50 \Omega}$

This is convenient.

What if there is no transmission line?

We can still use the chart, just define whatever Z_0 is convenient.

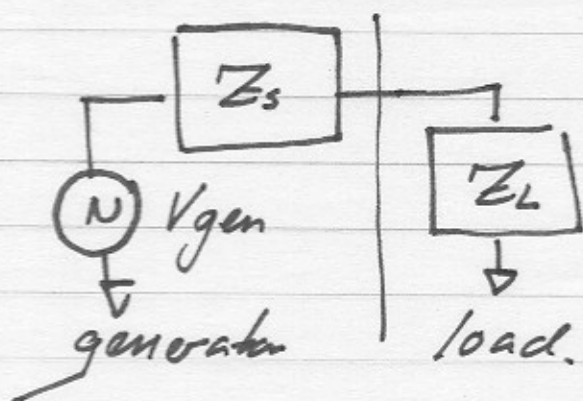


* solution on smith chart (next page)

gives $Z_{in} = 75 + j0\ \Omega$.

* This is also an introductory example of impedance matching.

First statement of impedance matching



$$Z_s = R_s + jX_s$$

Maximum power delivered to load if

$$Z_L = Z_s^* \quad (\text{conjugate matching})$$

$$P_{\max} = \|V_{\text{gen}}\|^2 / 4 \operatorname{Re}[Z_s]$$

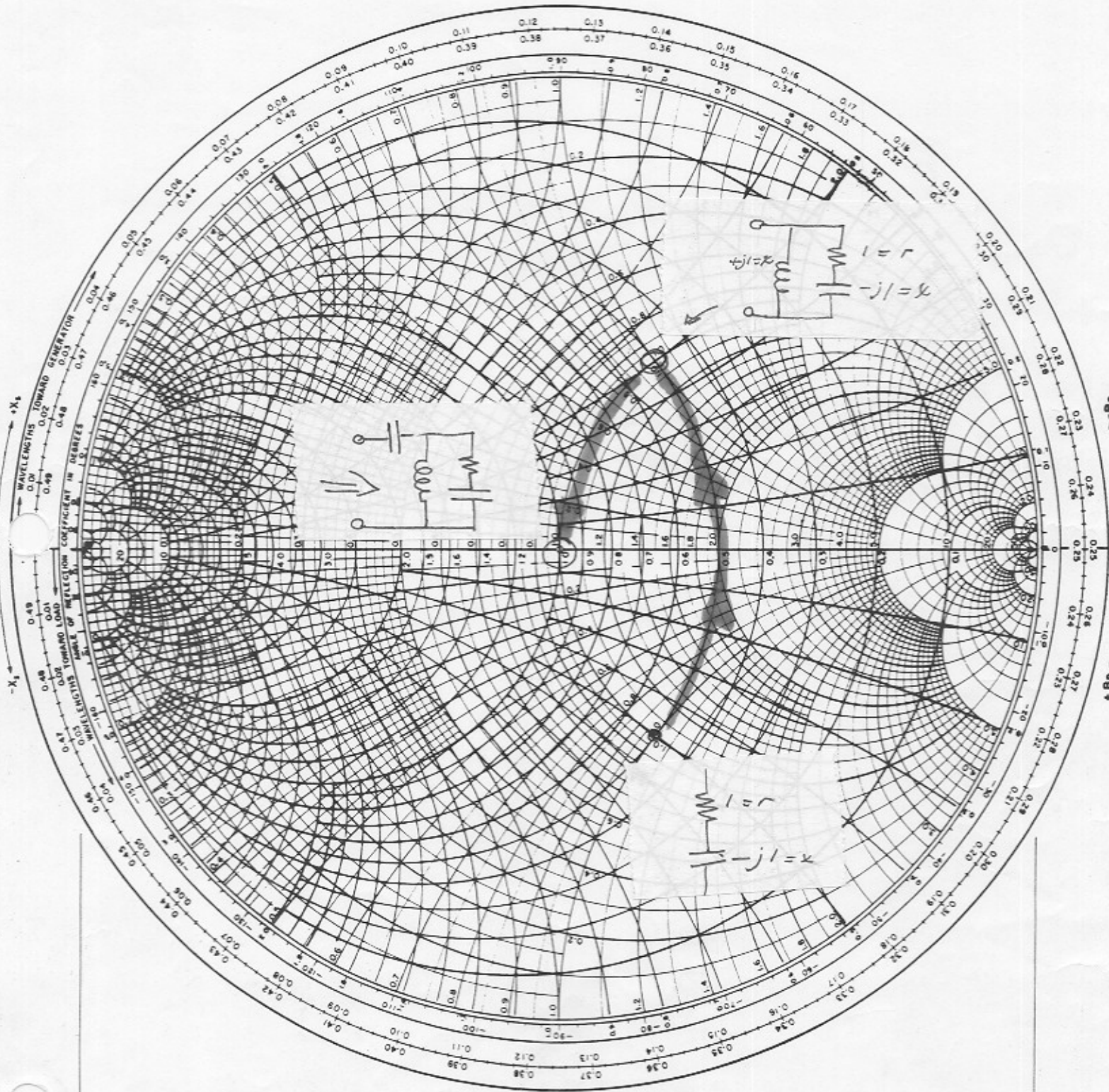
$$= \|V_{\text{gen}}\|^2 / 4 R_s$$

R.M.S. quantities are assumed.

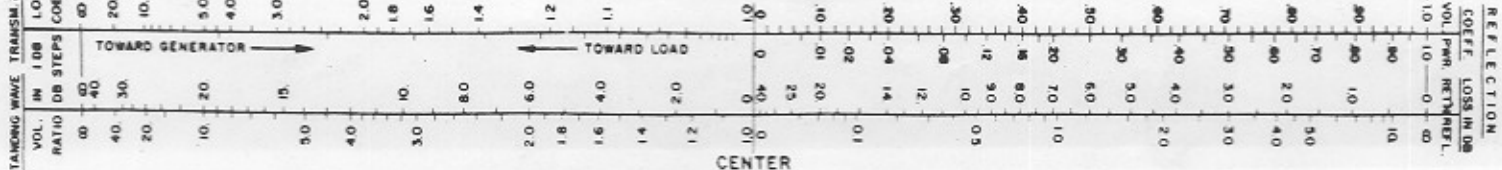
This is called the power available from the generator.

NAME	TITLE	DWG. NO.
SMITH CHART FORM ZY-01-N	ANALOG INSTRUMENTS COMPANY, NEW PROVIDENCE, N.J. 07974	DATE

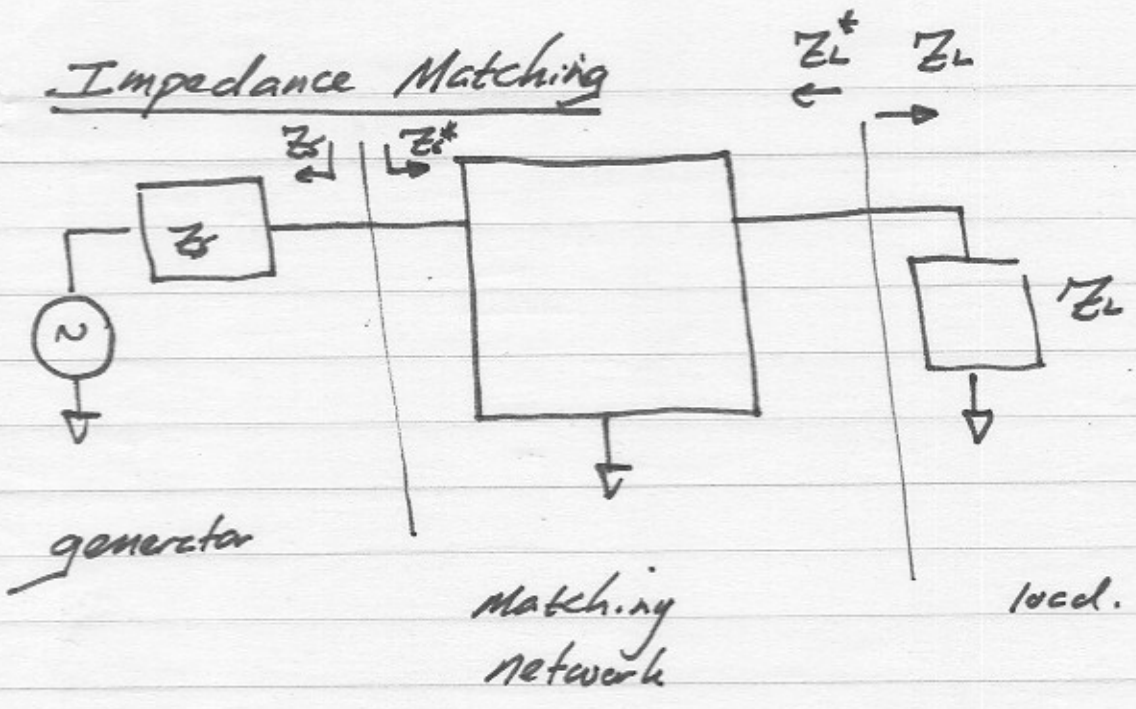
NORMALIZED IMPEDANCE AND ADMITTANCE COORDINATES



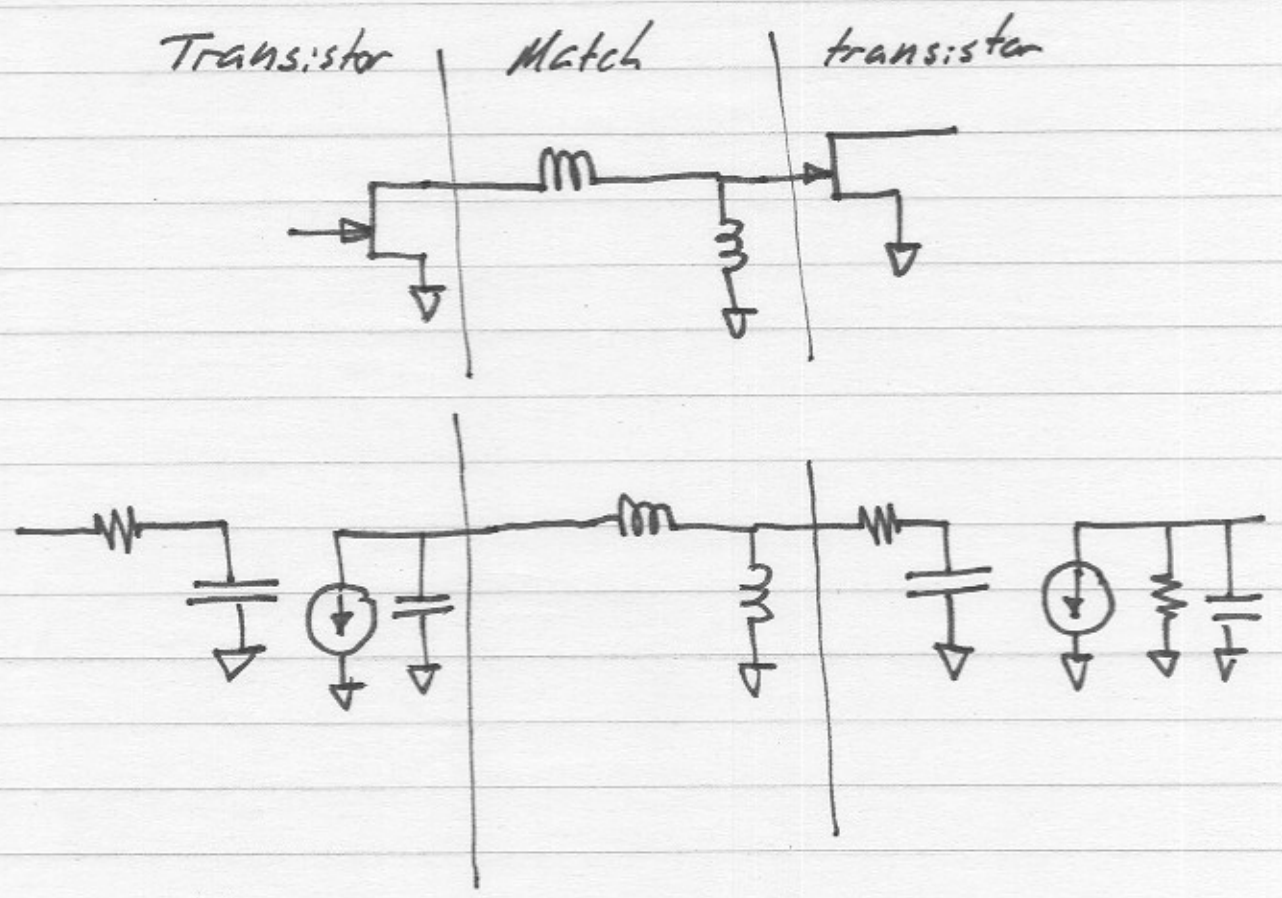
RADIALLY SCALED PARAMETERS



Impedance Matching



ex:



109 (37)
21

NAME	TITLE	DWG. NO.
SMITH CHART FORM ZY-01-N	ANALOG INSTRUMENTS COMPANY, NEW PROVIDENCE, N.J. 07974	DATE

NORMALIZED IMPEDANCE AND ADMITTANCE COORDINATES

