

ECE 202A Notes set 7

A review of circuit design methods

* Here we come to a decision point in the course: serve all students, or only those expert in circuit design?

* Let's try to serve all students, including those who are rusty in circuit design.

* The next six pages neatly summarize the things which you must review if unfamiliar with the topic.

ECE137D Notes set 6

Method of time constants:

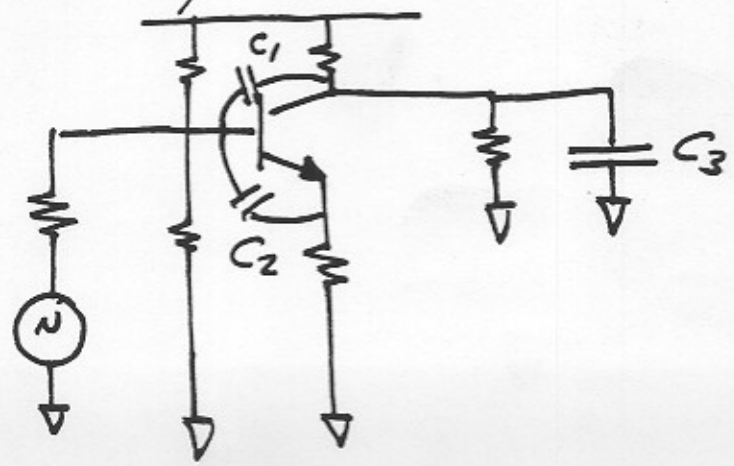
General Frequency response of circuit:

$$\frac{V_o(s)}{V_g(s)} = \frac{1 + b_1 s^2 + b_2 s^4 + \dots}{1 + a_1 s + a_2 s^2 + \dots} \cdot A_{dc}$$

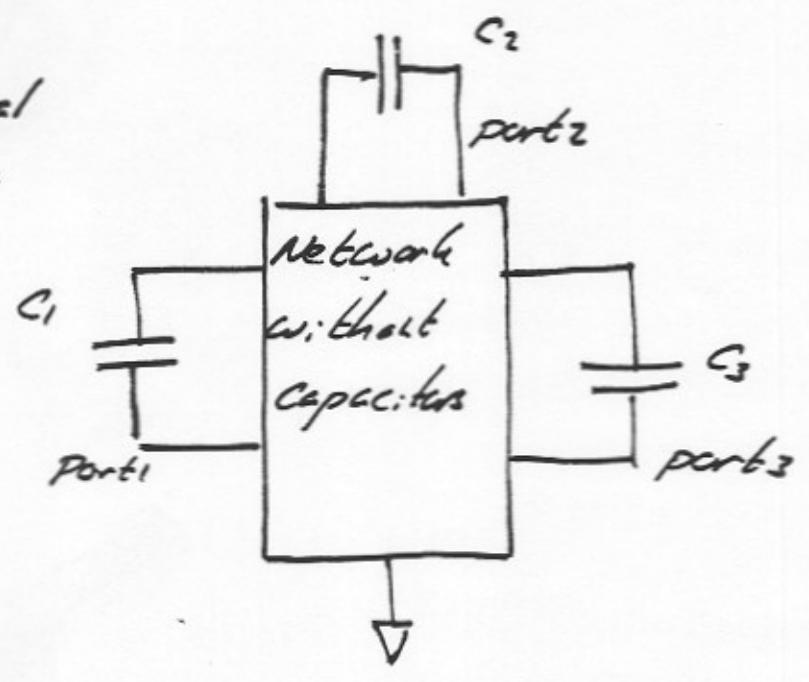
Methods of solution of circuit:

- Nodal Analysis: hard - and unrecognizable answers.
- Miller Approximations: often quite wrong.
- Time constant Method: accurate, fairly easy.

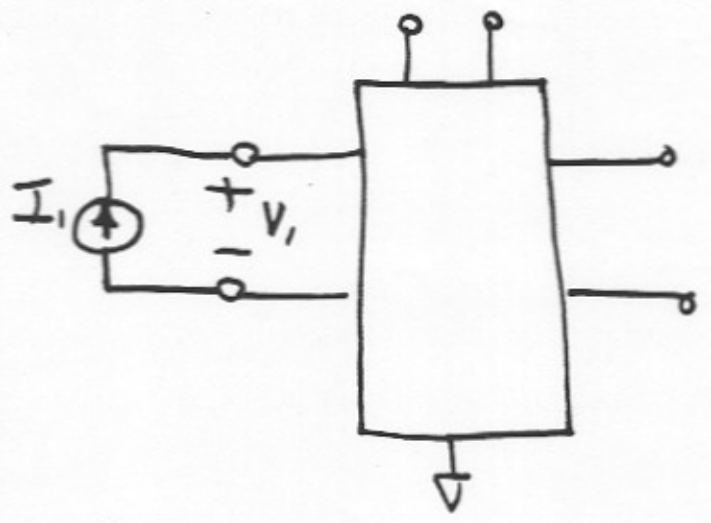
Network - An amplifier:



General
Picture



define R_{ii}^0 - Resistance at port i with all capacitors removed:

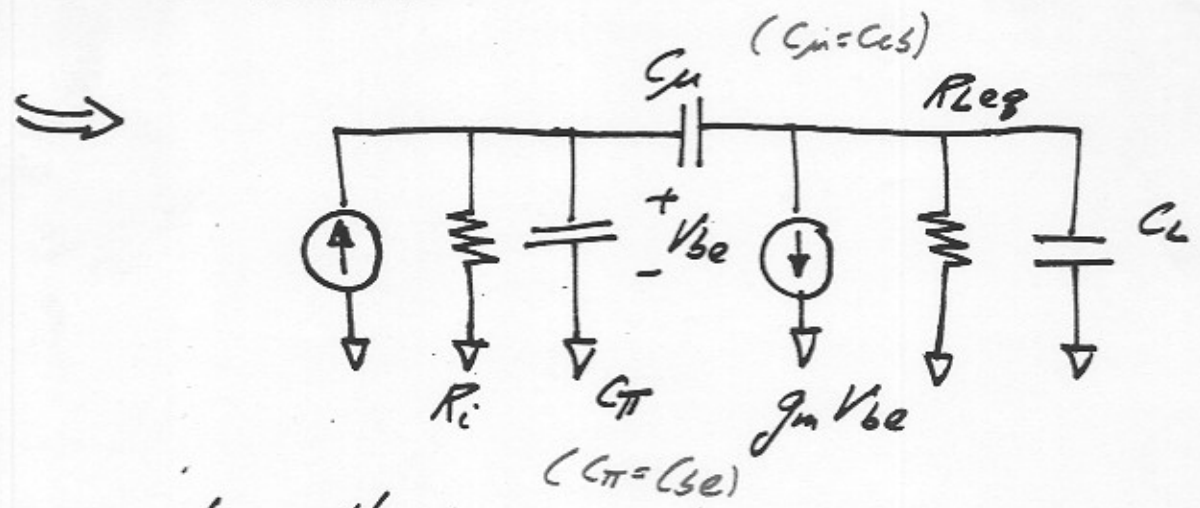
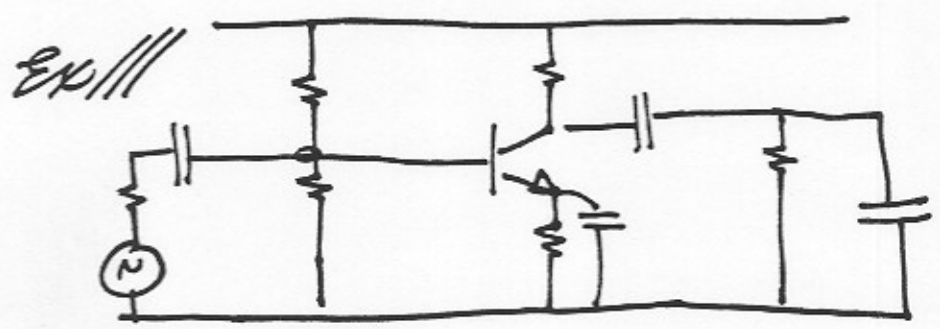


$$R_{ii}^0 \triangleq V / I$$

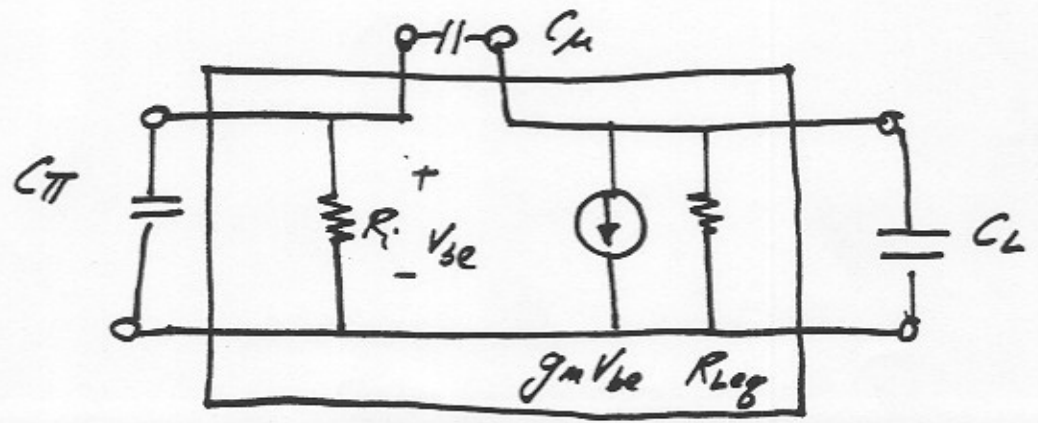
Method of time constants gives:

$$\tau_1 = R_{11}^0 C_1 + R_{22}^0 C_2 + R_{33}^0 C_3 + \dots$$

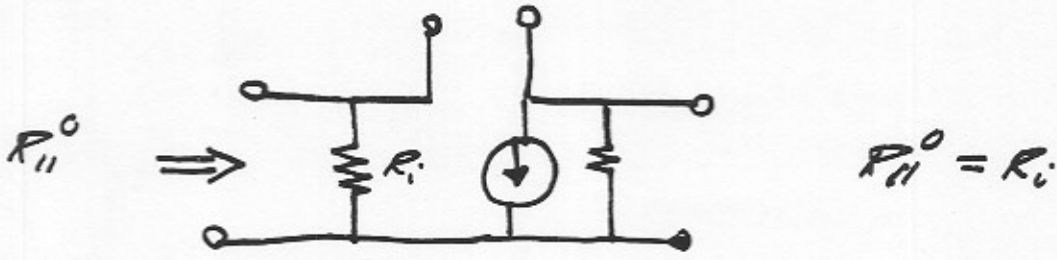
these are called the open-circuit time constants



Network without capacitors:

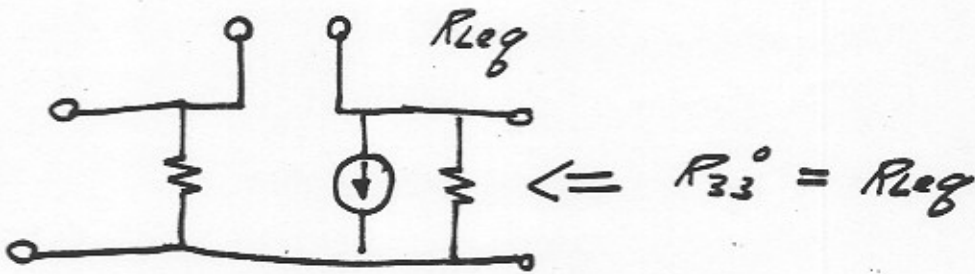


What is R_{11}^0 ? Easy!



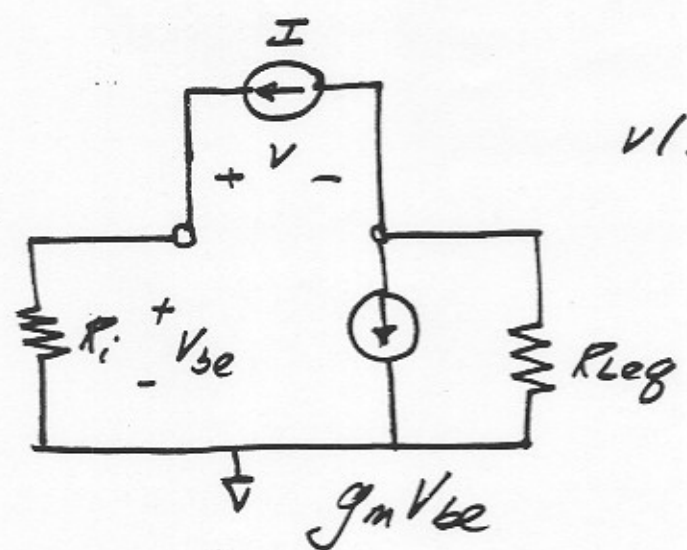
- so first time constant term is $C_{\pi} R_i$

What is R_{33}^0 ? Easy!

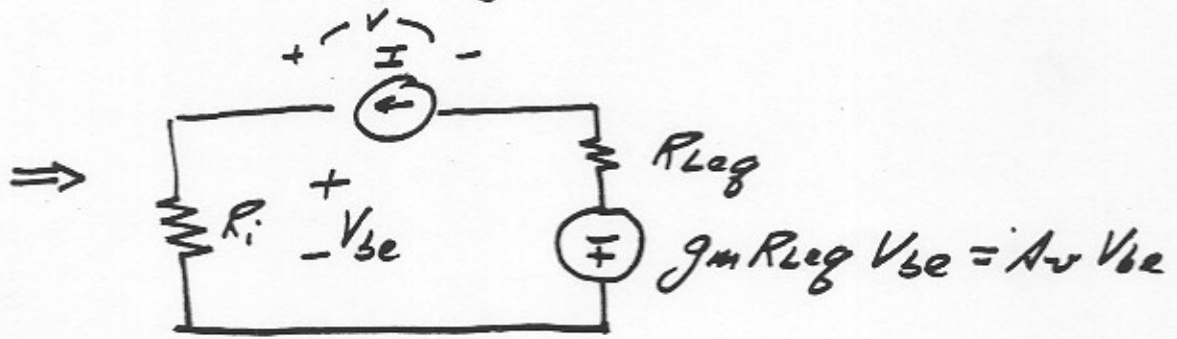


- so second time constant term is $C_L R_{leg}$.

What is R_{22}^0 ? - Harder



$V/I = R_{22}^0$



$$V = V_{be} + A_v V_{be} + I R_{leg}$$

$$= I R_i (1 + A_v) + I R_{leg}$$

$R_{22}^0 = R_i (1 + A_v) + R_{leg}$

-so third term in time constant is

$C_u [R_i (1 + A_v) + R_{leg}]$

Sum of time constants is

$$a_1 = R_i C_{\pi} + R_{leg} C_L + C_{\mu} [R_i (1 + A_v) + R_{leg}]$$

or grouping terms differently:

$$a_1 = R_i [C_{\pi} + C_{\mu} (1 + A_v)] + R_{leg} [C_L + C_{\mu}]$$

this we recognize as the same answer as given by both Nodal analysis & the Miller Effect.



What about higher-order poles?

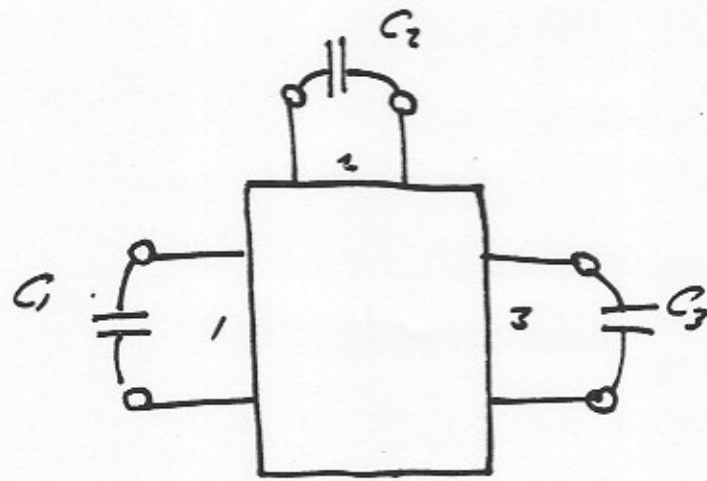
$$\frac{V_o(s)}{V_g(s)} = A_{MB} \frac{\text{Numerator (zeros)}}{1 + a_1 s + a_2 s^2 + a_3 s^3 + \dots}$$

If a_3 and higher terms are small, and: $\frac{a_2}{a_1} \ll a_1$, then:

$$\approx A_{MB} \frac{\text{Numerator}}{(1 + a_1 s)(1 + s \frac{a_2}{a_1})}$$

(separated-pole approximation)

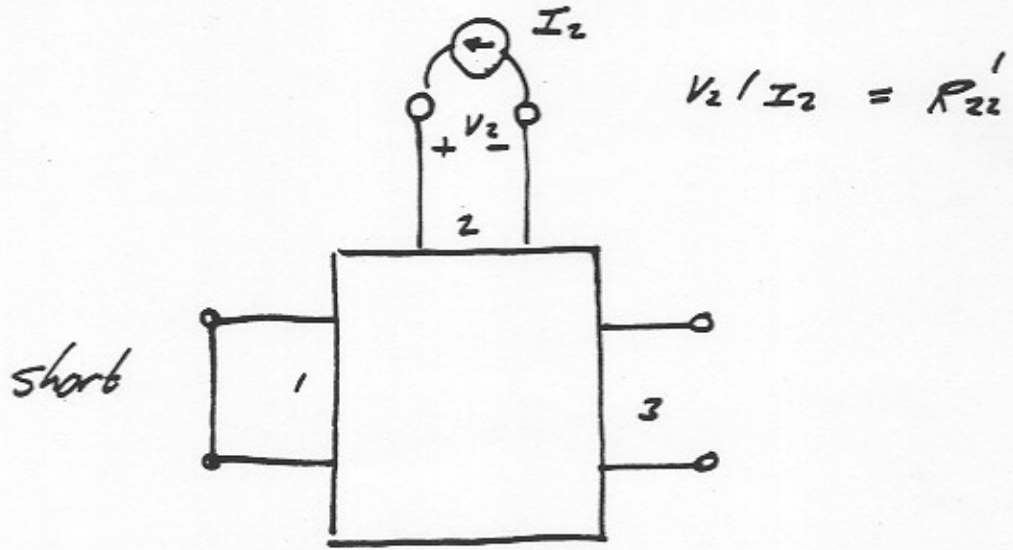
how do we find a_2 ?



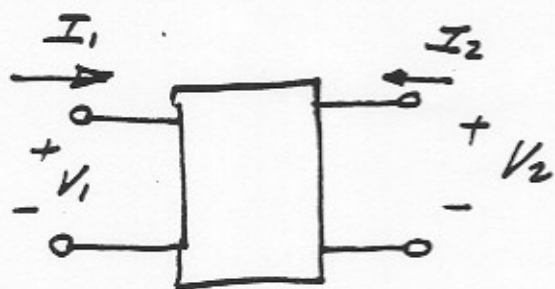
$$a_2 = R_{11}^0 C_1 R_{22}^1 C_2 + R_{11}^0 C_1 R_{33}^1 C_3 + R_{22}^0 C_2 R_{33}^2 C_3$$

where we define R_{22}^1 as below:

"Resistance seen at port 2 with port 1 shorted"



Now Consider a 2-port For a Moment:



$$\begin{bmatrix} V_1 \\ V_2 \end{bmatrix} = \begin{bmatrix} Z_{11} & Z_{12} \\ Z_{21} & Z_{22} \end{bmatrix} \begin{bmatrix} I_1 \\ I_2 \end{bmatrix}$$

$$R_{11}^0 = Z_{11} = \frac{V_1}{I_1} \Big|_{I_2=0} \quad \& \quad R_{11}^2 = \frac{V_1}{I_1} \Big|_{V_2=0}$$

$$\text{If } V_2 = 0 \Rightarrow Z_{21} I_1 + Z_{22} I_2 = 0 \Rightarrow \frac{V_1}{I_1} = \boxed{Z_{11} + \frac{Z_{12}(-Z_{21})}{Z_{22}} = R_{11}^2}$$

$$\text{Similarly: } R_{22}^1 = Z_{22} - \frac{Z_{12}Z_{21}}{Z_{11}}$$

$$\text{So: } R_{11}^0 R_{22}^1 = Z_{11} \left[Z_{22} - \frac{Z_{12}Z_{21}}{Z_{11}} \right] = Z_{11}Z_{22} - Z_{12}Z_{21} = \Delta Z$$

$$\text{and } R_{11}^2 R_{22}^0 = \left[Z_{11} - \frac{Z_{12}Z_{21}}{Z_{22}} \right] Z_{22} = \dots = \Delta Z$$

$$\underline{\text{So:}} \quad \boxed{R_{11}^0 R_{22}^1 = R_{22}^0 R_{11}^2 = \Delta Z}$$

So

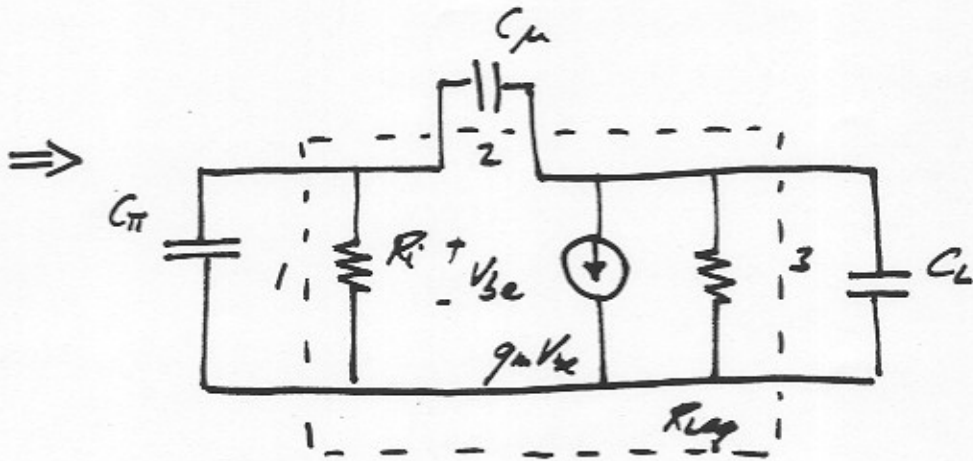
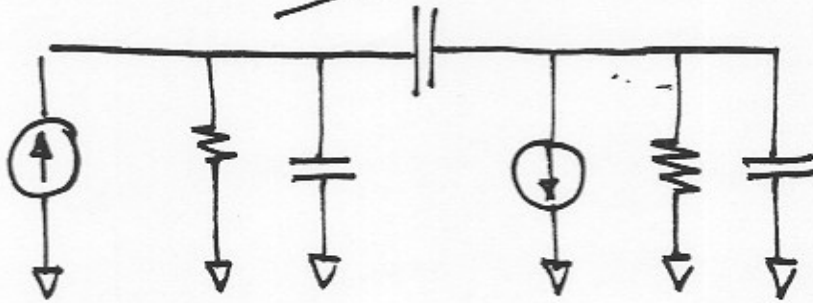
$$a_2 = R_{11}^0 R_{22}^1 C_1 C_2 + R_{11}^0 R_{33}^1 C_1 C_3 + R_{22}^0 R_{33}^2 C_2 C_3$$
$$= R_{22}^0 R_{11}^2 C_1 C_2 + R_{33}^0 R_{11}^3 C_1 C_3 + R_{33}^0 R_{22}^3 C_2 C_3$$

each term can be calculated 2 ways



ex III

Common-emitter stage



$R_{11}^o R_{22}^i$:

$$R_{11}^o = R_i, \quad R_{22}^i = R_{Leg}$$

$$\underline{R_{11}^0 R_{33}^1} \quad R_{11}^0 = R_i \ ; \ R_{33}^1 = R_{Leg}$$

$$\underline{R_{22}^0 R_{33}^2} = R_{33}^0 R_{22}^3$$

$$R_{33}^0 = R_{Leg} \ ; \ R_{22}^3 = R_i$$

so: $R_{11}^0 R_{22}^1 = R_{11}^0 R_{33}^1 = R_{22}^0 R_{33}^2 = R_i R_{Leg}$.

$$\Rightarrow a_2 = [C_{\mu} C_L + C_{\mu} C_{\pi} + C_{\pi} C_L] R_i R_{Leg}$$

So:

$$\frac{V_o(s)}{V_g(s)} = A_{mb} \frac{\text{Numerator}}{1 + a_1 s + a_2 s^2}$$

$$a_1 = R_i [C_{\pi} + C_{\mu}(1 + g_m R_{log})] + R_{log} [C_{\mu} + C_L]$$

$$a_2 = R_i R_{log} [C_{\mu} C_{\pi} + C_{\mu} C_L + C_{\pi} C_L]$$

• same answer as by Nodal Analysis.

Method of Time Constants

	$\frac{V_{out}}{V_{gen}} = \left(\frac{V_{out}}{V_{gen}} \right)_{mb} \frac{1 + b_1 s + b_2 s^2 + \dots}{1 + a_1 s + a_2 s^2 + \dots}$ $a_1 = R_{11}^0 C_1 + R_{22}^0 C_2 + R_{33}^0 C_3 + R_{44}^0 C_4$ $a_2 = R_{11}^0 R_{22}^1 C_1 C_2 + R_{11}^0 R_{33}^1 C_1 C_3 + R_{11}^0 R_{44}^1 C_1 C_4 + R_{22}^0 R_{33}^2 C_2 C_3 + R_{22}^0 R_{44}^2 C_2 C_4 + R_{33}^0 R_{44}^3 C_3 C_4$ <p>notethat $R_{xx}^0 R_{yy}^x = R_{xx}^y R_{yy}^0$</p>
	$R_{xx}^0 = R_x \left(r_e \parallel R_{Leq} + R_i (1 - A_{vmb}) \right)$ $A_{vmb} = \left(R_{Leq} / (r_e + R_{Leq}) \right)$
	$R_i = R_x \parallel R_x$ $R_{yy}^0 = R_i (1 + g_m R_{Leq}) + R_e + R_{Leq}$
	<p>For the bipolar</p> $R_{in,base} = (\beta + 1)(R_E + r_e)$ $R_{in,emitter} = \left(r_e + \frac{R_B}{\beta + 1} \right) \left(\frac{r_o + R_C}{r_o} \right) \cong \left(r_e + \frac{R_B}{\beta + 1} \right)$ $R_{in,collector} = r_o \left\{ 1 + g_m R_E \left(\frac{r_o}{r_o + R_E + R_B} \right) \right\} \cong r_o \{ 1 + g_m R_E \}$ <p>these give almost all the remaining relationships needed for the MOTC</p>

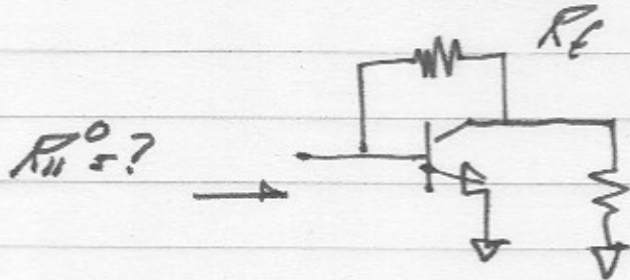
For the fet: THE EQUATIONS ARE THE SAME IF B IS INFINITE!

If we refer to the MOTC crib sheet on the previous page, we will find expressions for R_{xx}^y & R_{xx}^o for almost every case

It is generally unnecessary with the MOTC to do mesh mth!

The only exception:

networks with feedback, e.g.

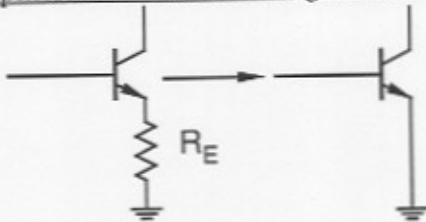


will generally need working through on a case-by-case basis.

Emitter Degeneration

From the crib sheet:

Simplification of Emitter Degeneration



approximate only; check notes for bounds on validity

$$\tilde{C}_x = C_x (r_e / (r_e + R_E))$$

$$\tilde{g}_m = g_m (r_e / (r_e + R_E)) = \alpha / (r_e + R_E)$$

$$\tilde{C}_\mu = C_\mu$$

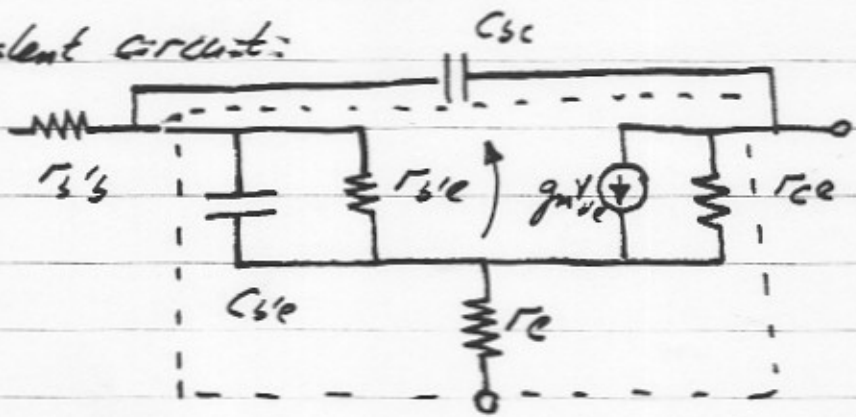
$$\tilde{R}_x = R_x ((r_e + R_E) / r_e) = (\beta + 1)(r_e + R_E)$$

this is a powerful simplification for the purposes of analysis.

Proof (and limits) of the relationship are given on the following pages, which will not be covered in lecture...

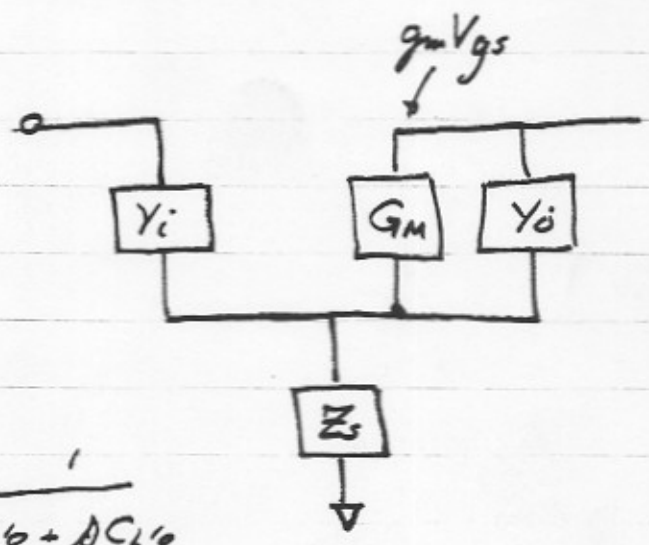
conversion (first-order) of device equivalent circuit to 'absorb' R_s (or R_E)

equivalent circuit:



First: put aside r_s 's & $C_s'c$; we will put back later

Second: some notation:



$$Z_i = \frac{1}{Y_i} = \frac{1}{r_s'e + \omega C_s'e}$$

$$Y_o = \frac{1}{Z_o} = \frac{1}{r_c'e} \quad (\text{not characteristic impedance})$$

$$Z_s = r_e = 1/Y_s : \text{tendency to mix notation: } r_s = r_e$$

Approximations

We are developing only an approximate equivalent circuit.

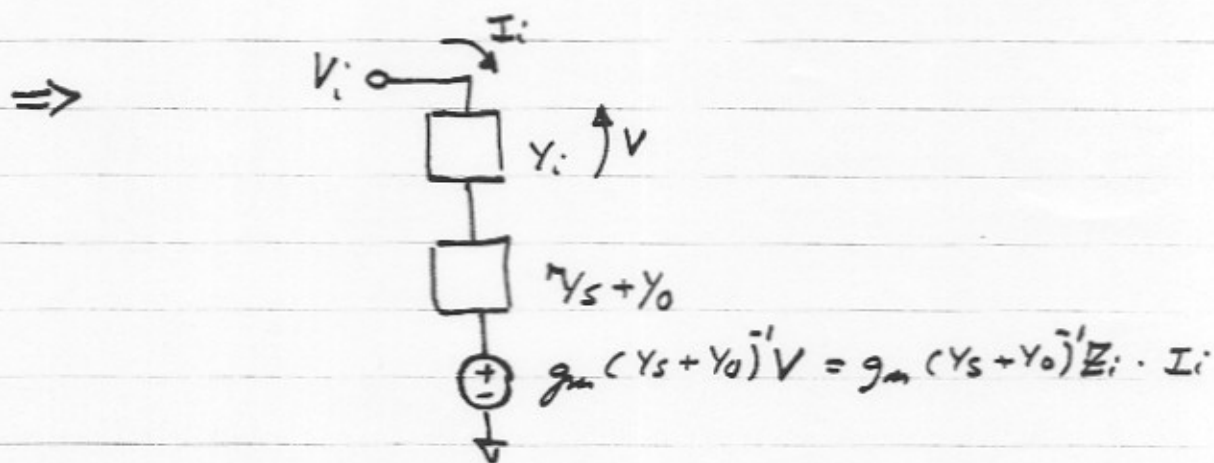
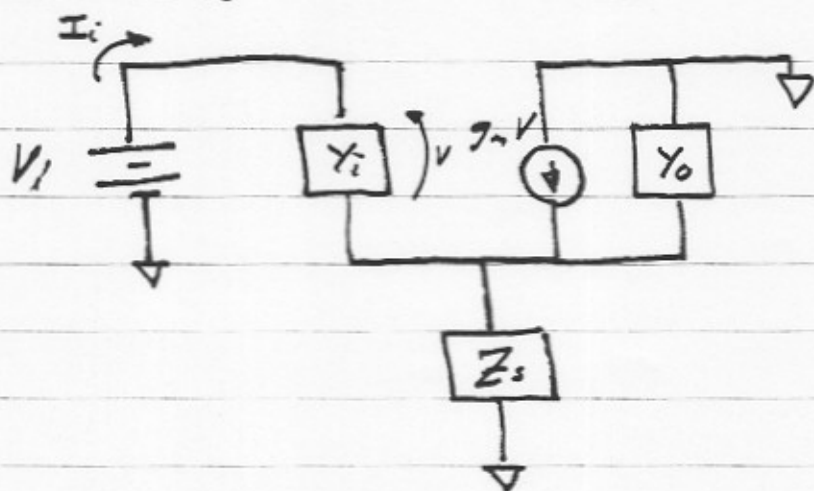
Assuming: $Z_s \ll Z_i$ otherwise, expressions can be generalized

$g_m \gg Y_i$ this implies $f \ll f_{T1}$

$Z_s \ll \frac{1}{Y_0}$ reasonable almost always.

Get the Y-parameters:

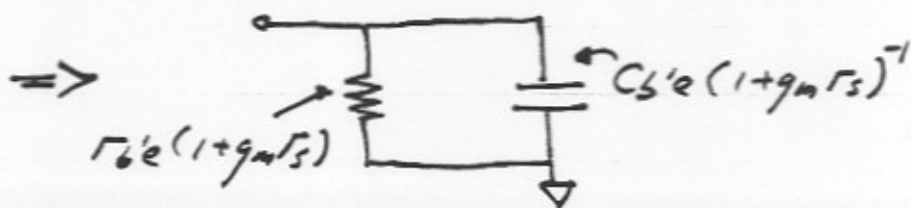
$$Y_{11} = I_i / V_i \mid V_o = 0$$



$$\frac{1}{Y_{11}} = Z_i + \frac{1}{Y_s + Y_o} + Z_i \frac{g_m}{Y_s + Y_o} \cong Z_i (1 + g_m Z_s) + Z_s$$

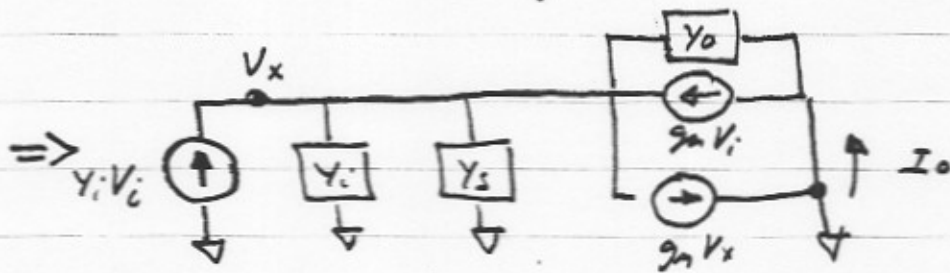
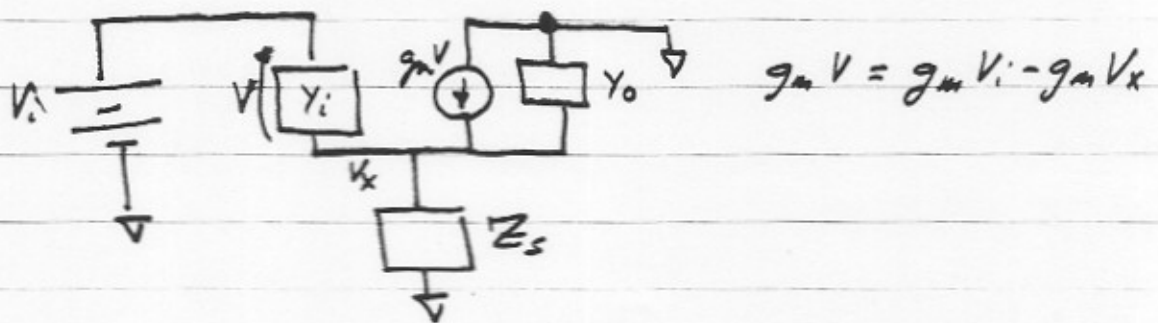
$$\cong Z_i (1 + g_m Z_s)$$

$$Y_{11} \cong \frac{Y_i}{1 + g_m r_s}$$



Transconductance

$$Y_{21} = I_o / V_i \big|_{V_o = 0}$$



$$V_x (Y_i + Y_s + Y_o + g_m) - (Y_i + g_m) V_i = 0$$

$$V_x = V_i (Y_i + g_m) / (Y_i + Y_s + Y_o + g_m)$$

$$I_o = -(g_m + Y_o) V_x + g_m V_i$$

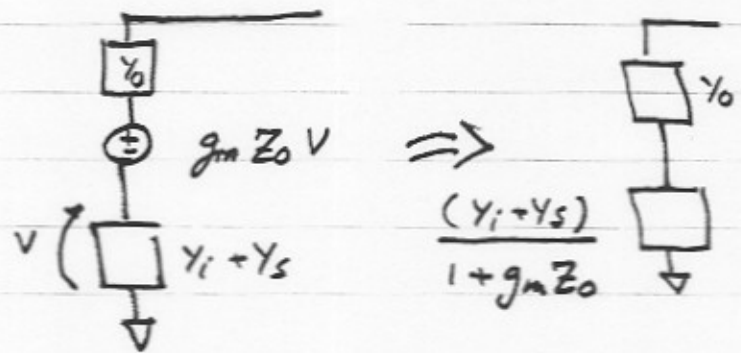
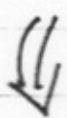
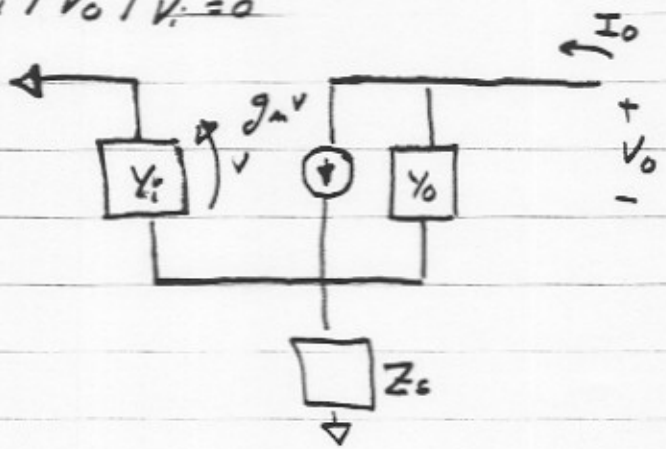
$$I_o / V_i = \frac{-Y_o Y_i + g_m Y_s}{g_m + Y_o + Y_i + Y_s}$$

$$\approx \frac{g_m Y_s}{g_m + Y_s} = \frac{g_m}{1 + g_m Y_s}$$

$$\Rightarrow \downarrow g_m (1 + g_m Y_s)^{-1}$$

output conductance

$$Y_{22} = I_i / V_o \mid V_i = 0$$



$$Z_{out} = \frac{1}{Y_{22}} = Z_o + (1 + g_m Z_o) / (Y_i + Y_s)$$

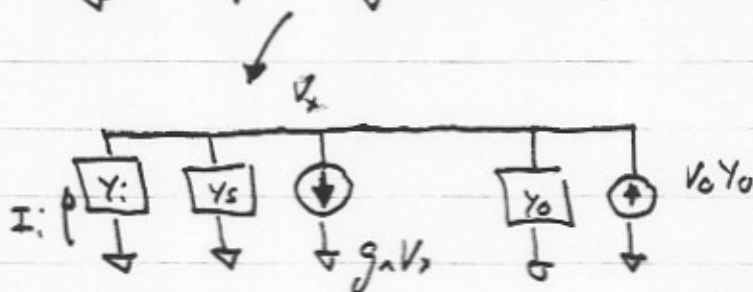
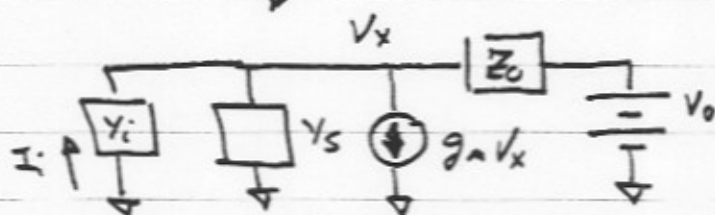
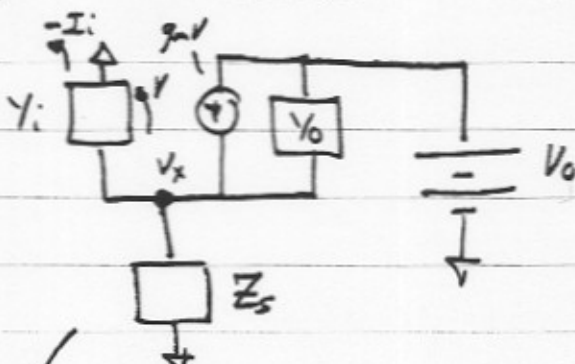
$$\cong Z_o + (1 + g_m Z_o) / Y_s = Z_o + (1 + g_m Z_o) r_s$$

$$= Z_o (1 + g_m r_s) + r_s \cong Z_o (1 + g_m r_s)$$

$$\cong r_{ce} \cdot (1 + g_m r_s)$$

Reverse Conductance: (Y_{12})

$$Y_{12} = \frac{I_i}{V_o} \Big|_{V_i=0}$$



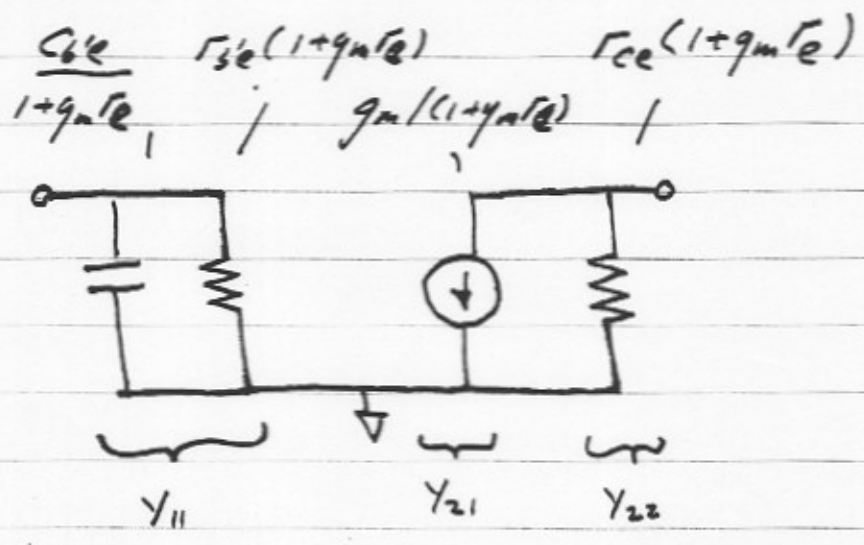
$$V_x (Y_i + Y_s + g_m + Y_o) = V_o Y_o$$

$$I_i = -Y_i V_x$$

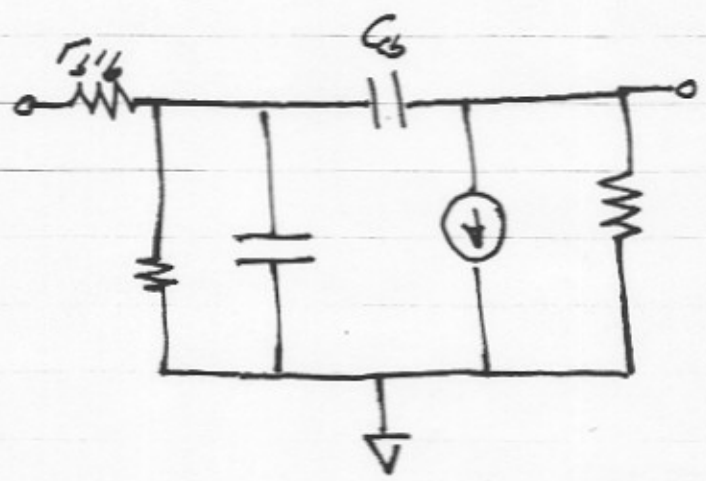
$$\frac{I_i}{V_o} = \frac{-Y_i Y_o}{Y_i + Y_s + g_m + Y_o} \approx -Y_i \frac{Y_o}{g_m + Y_s}$$

but Y_i is fairly small and $Y_o \ll g_m \Rightarrow$ neglect to first order.

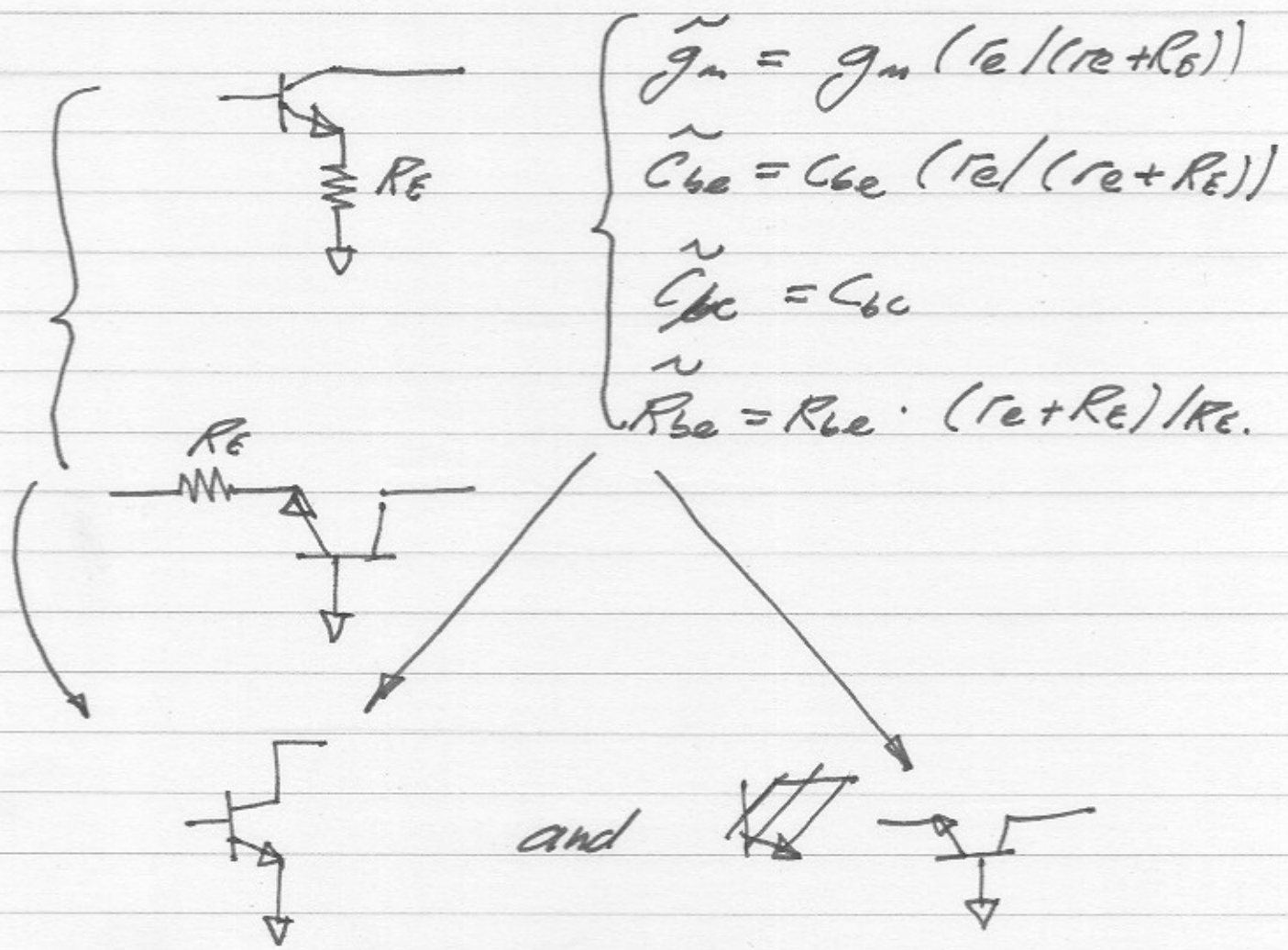
Bringing it all together:



Adding back parasitic C_{cb} , $r_{s'b}$:

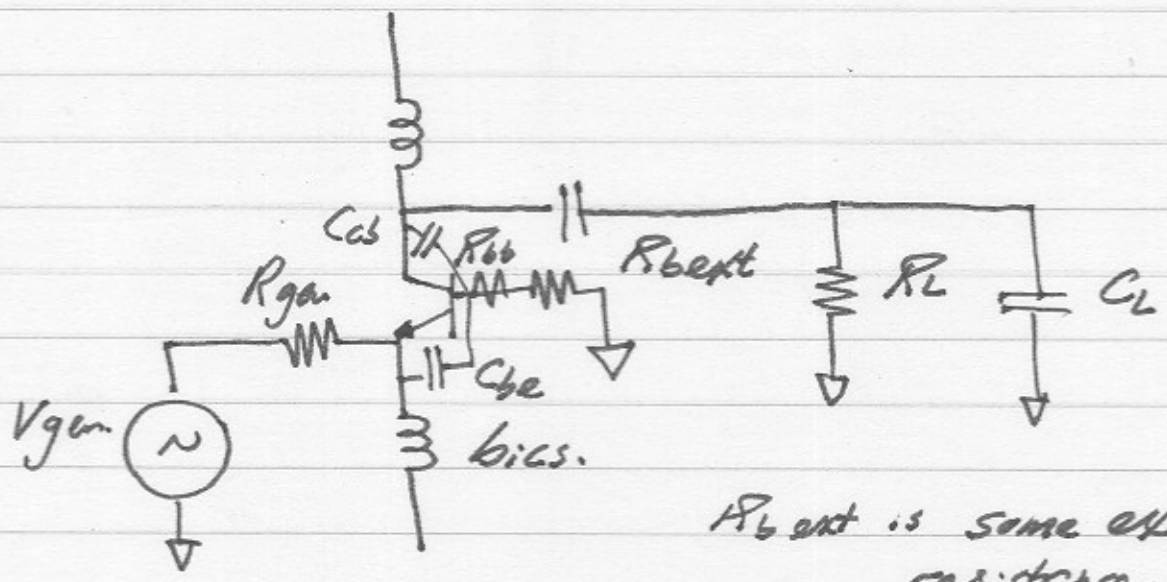


Note that, within the bounds of accuracy of the derivation above, the degeneration method can be used on both common-emitter and common-base stages:



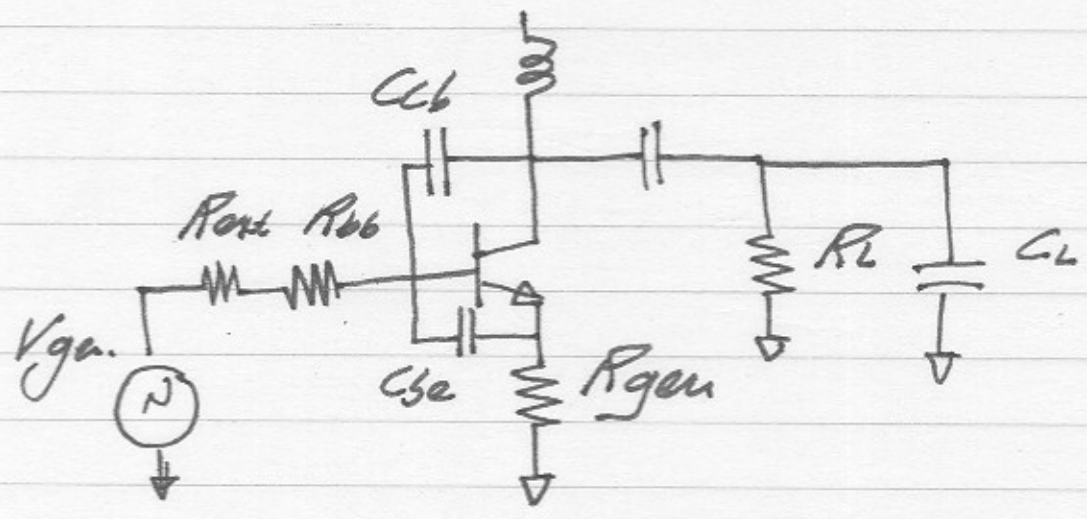
Common-base stage by MOTC:

now lets include a non-zero R_{bb} !



$R_{b\ ext}$ is some external resistance...

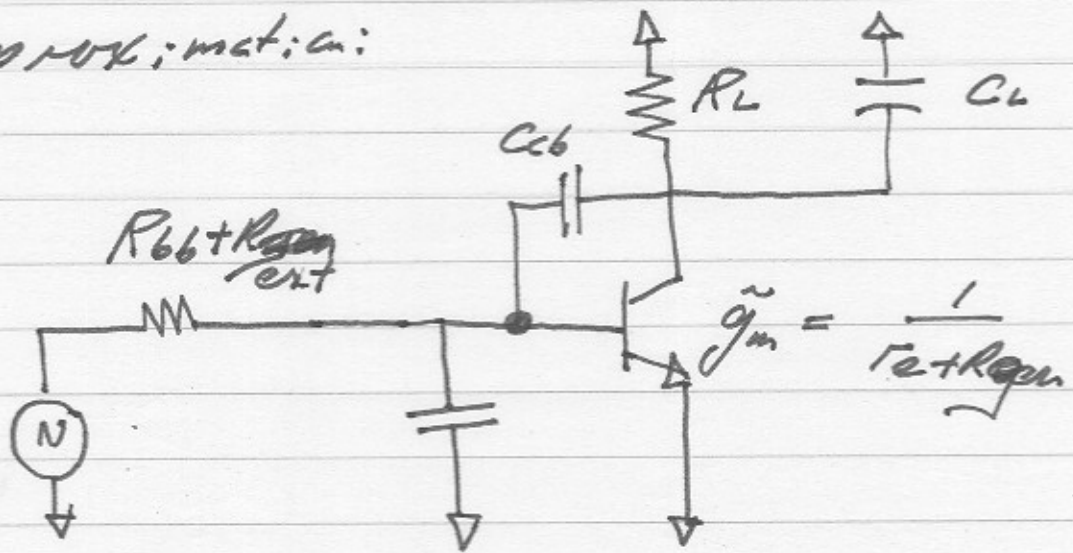
Before taking this further, consider a common-emitter circuit with peculiar notation:



* Lets first simplify slightly by taking $\beta \rightarrow \infty$.

* Then lets use the degeneration approximation:

approximation:



$$\tilde{C}_{be} = C_{be} \cdot \frac{r_e}{r_e + R_{gen}}$$

Now we know that

(From nodal analysis or MOTC)

$$H(A) \propto \frac{N(A)}{1 + a_1 A + a_2 A^2}$$

where

$$a_1 = (R_{bb} + R_{ext}) \left[\frac{C_{be} r_e}{r_e + R_{gen}} + C_{cb} \left(1 + \frac{R_L}{r_e + R_{gen}} \right) \right]$$

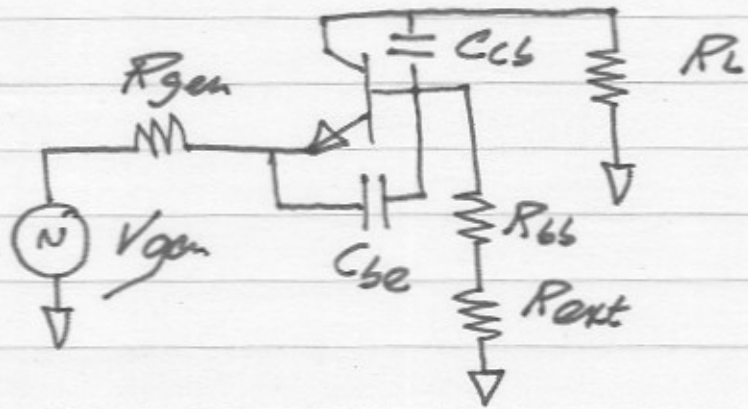
$$a_2 = C_{cb} \cdot R_L + C_L \cdot R_L$$

$$a_2 = R_L \cdot (R_{bb} + R_{ext})$$

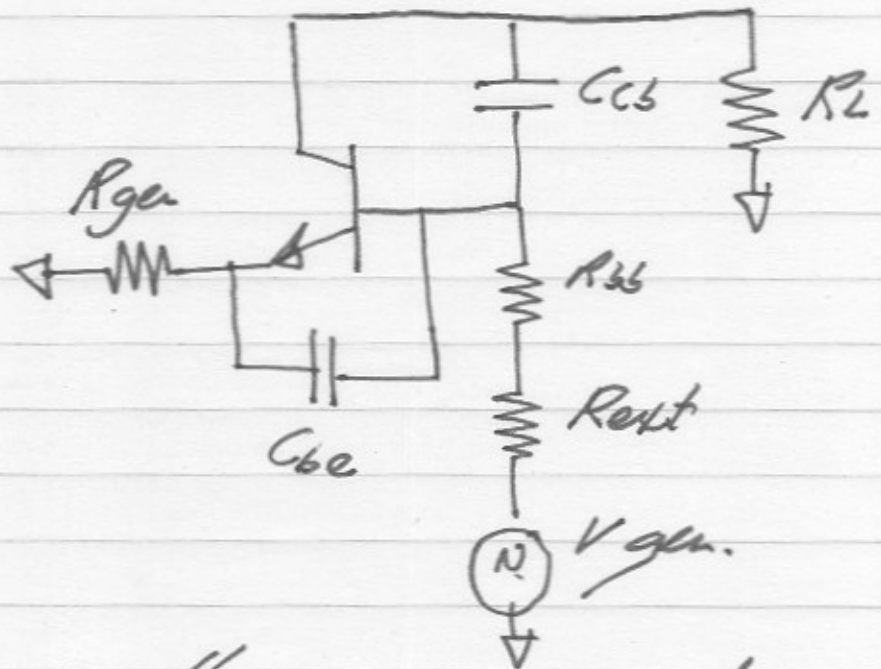
$$\cdot \left[C_{cb} \cdot C_L + \frac{C_{be} r_e}{r_e + R_{gen}} \cdot C_L + \frac{C_{be} r_e}{r_e + R_{gen}} \cdot C_{cb} \right]$$

But wait:

CB circuit:

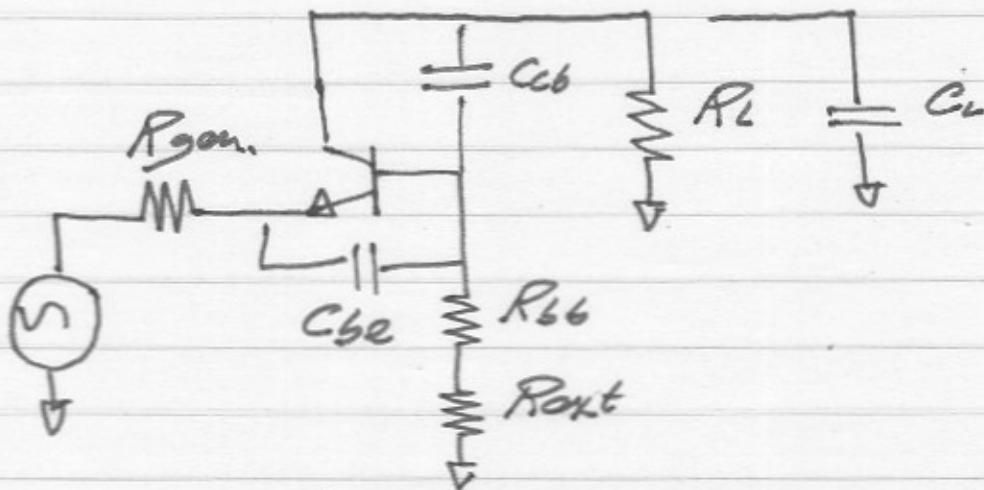


CE circuit



the 2 circuits differ only in the connection of V_{gen} , and that has no effect on a_{11} or a_{22} or a_{33} etc (only b_1 , b_2 , b_3)

So for the common base stage:



$$a_1 \approx C_{be} \cdot \left[\frac{r_e}{r_e + R_{gen}} \right] \cdot (r_{bb} + R_{ext})$$

$$+ C_{cb} \left[R_L + (R_{bb} + R_{ext}) \left(1 + \frac{R_L}{r_e + R_{gen}} \right) \right]$$

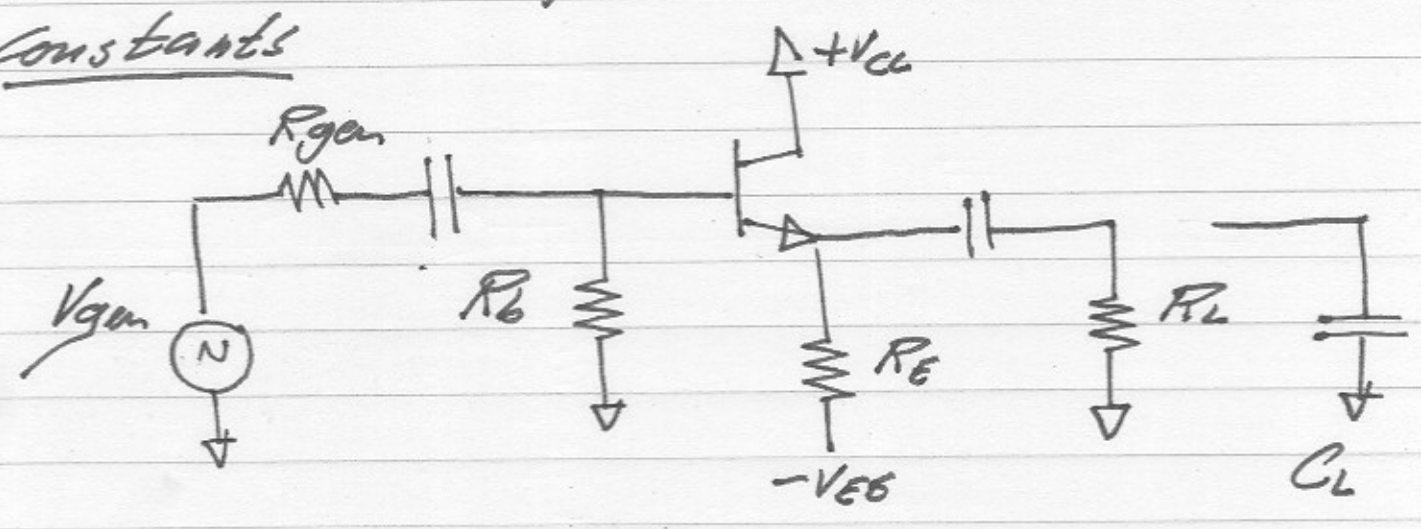
$$+ C_L \cdot R_L$$

this is very similar to the

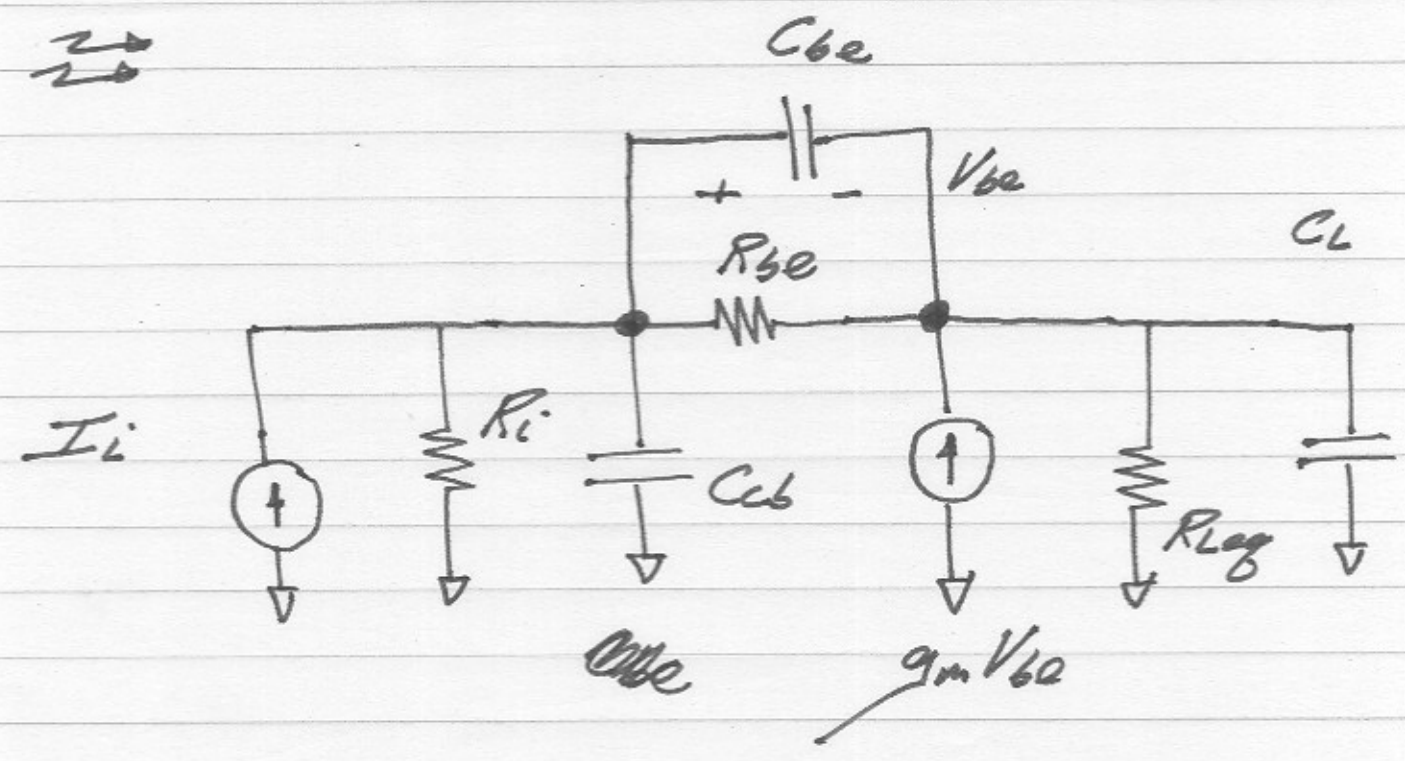
Miller Effect. $(R_{bb} + R_{ext})$ needs to be small.

Emitter follower by Method of time constants

Constants



↔

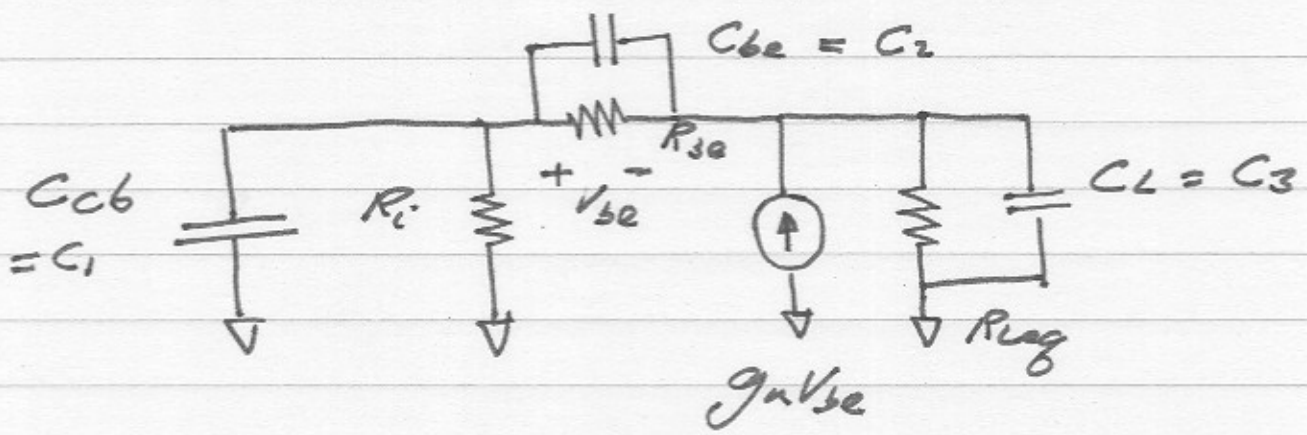


$$R_i = R_{B'} + (R_{gen} \parallel R_i)$$

$$R_{eq} = R_E \parallel R_L$$

Method of time constants:

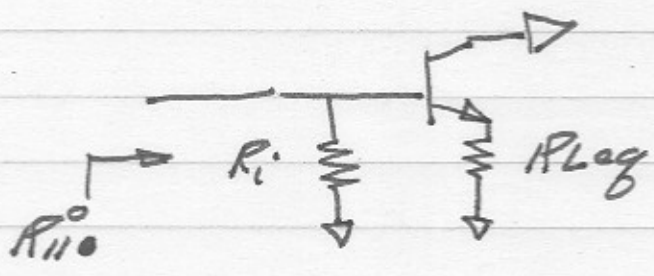
$$H(s) = H_{MB} \cdot \frac{N(s)}{1 + a_1 s + a_2 s^2}$$



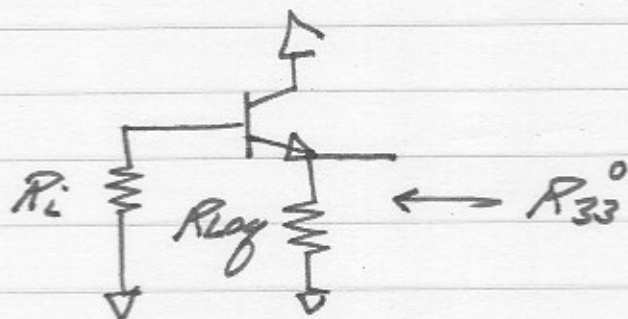
$$a_1 = C_1 R_{11}^0 + C_2 R_{22}^0 + C_3 R_{33}^0$$

$$\frac{C_1 R_{11}^0}{}$$

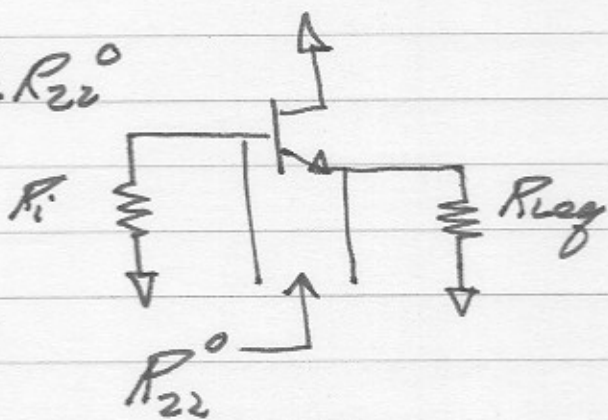
What is R_{11}^0 ?



by inspection $R_{11}^0 = R_i \parallel (\beta + 1)(r_e + R_{leg})$!

C₃ R₃₃^o

by inspection, $R_{33}^o = R_{leg} \parallel \left(r_e + \frac{R_i}{\beta + 1} \right)$

C₂ R₂₂^o

this is harder. Requires nodal analysis.

Answer is on crib sheet.

$$R_{22}^o = R_{\pi} \parallel \left[(r_e \parallel R_{leg}) + R_i \left(1 - \frac{R_{leg}}{r_e + R_{leg}} \right) \right]$$

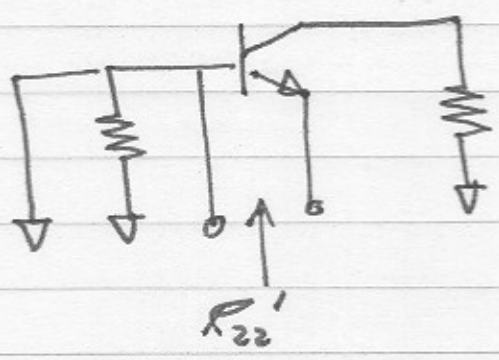
the very alert reader will recognize the Miller effect in the last term.

Now second-order time constants.

$$\tau_2 = C_1 C_2 R_{11}^0 R_{22}^1 + C_1 C_3 R_{11}^0 R_{33}^1 + C_2 C_3 R_{22}^0 R_{33}^2$$

R_{11}^0, R_{22}^0 & R_{33}^0 we have already found.

$R_{22}^1 = ?$

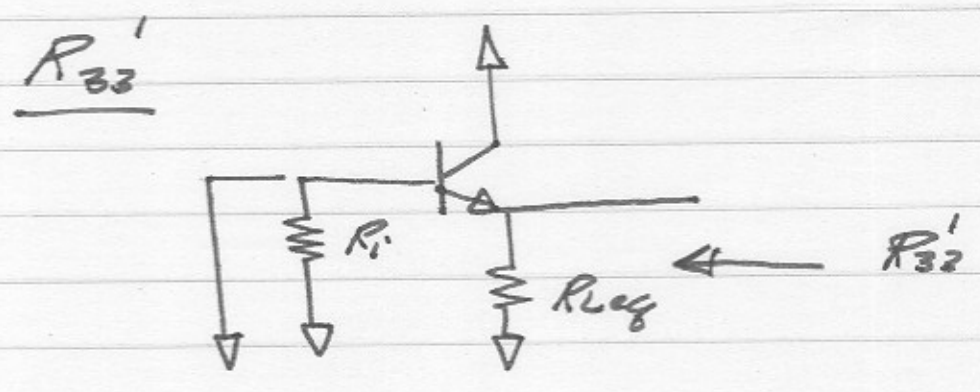


$$R_{22}^0 = R_{\pi} \parallel \left[(r_e \parallel R_{sig}) + R_i \left(1 - \frac{\beta_{sig}}{\beta_0 + \beta_{sig}} \right) \right]$$

we have just shorted part 1 (R_i)

so just set $R_i \rightarrow 0$.

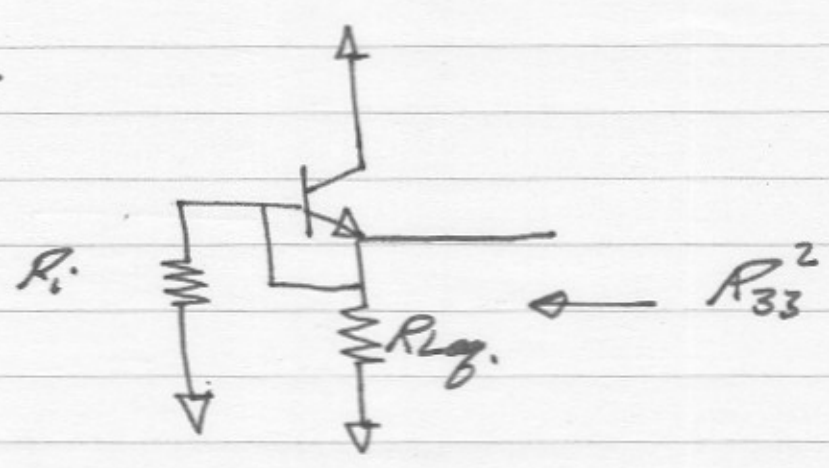
$$R_{22}^1 = R_{\pi} \parallel r_e \parallel R_{sig} \approx r_e \parallel R_{sig}$$



by same argument:

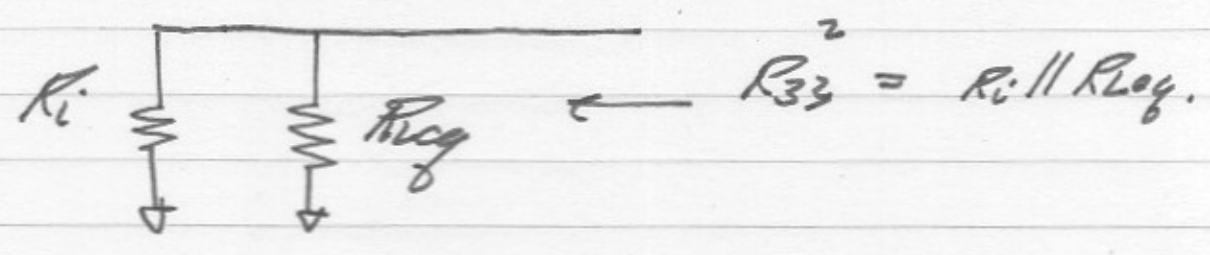
$$R_{33}' = R_{e2} \parallel \left(r_e + \frac{R_i}{\beta + 1} \right)$$
$$= R_{e2} \parallel r_e$$

R_{33}^2 :



b.e. junction is ac shorted.

no small-signal base or emitter current flows \rightarrow



Gathering terms:

$$a_1 = C_{be} \cdot \left[R_{be} \parallel \left(r_e \parallel R_{eq} + R_i \left(1 - \frac{R_{eq}}{R_{eq} + R_e} \right) \right) \right]$$

$$+ C_{cb} \left[R_i \parallel (\beta + 1)(r_e + R_{eq}) \right]$$

$$+ C_L \left[R_{eq} \parallel \left(r_e + \frac{R_i}{\beta + 1} \right) \right]$$

$$a_2 = (r_e \parallel R_{eq}) (R_i \parallel (\beta + 1)(r_e + R_{eq}))$$

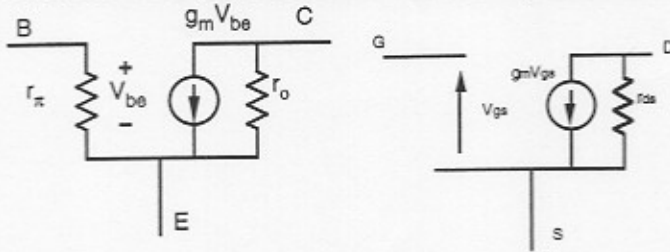
$$\cdot [C_{be} C_{cb} + C_{be} C_L + C_{cb} C_L]$$

\rightarrow and the bandwidth is

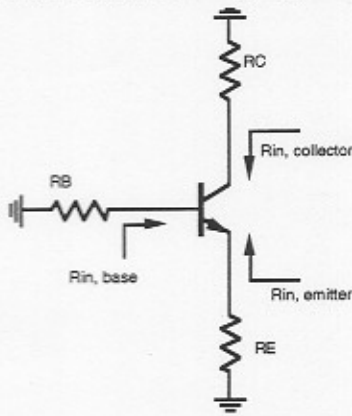
$$f_{p1} \sim \frac{1}{2\pi a_1}$$

... note there is also a zero @ $f_z = \frac{g_m}{2\pi C_{be}} = f_T$

Basics: Amplifiers at Low Frequencies



Left is the equivalent circuit of a bipolar transistor; $g_m = I_E / V_T = 1/r_e$, $V_T = kT/q$, $r_{\pi} = (\beta + 1)r_e = \beta/g_m$. On the right is the FET model. Below is a transistor with ac equivalent load resistances R_C , R_E , and R_B .

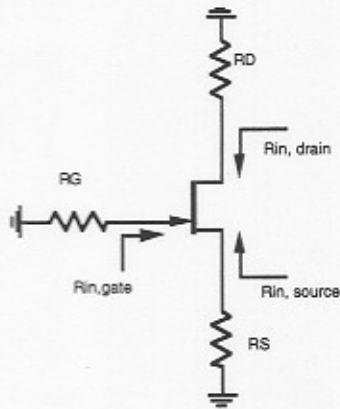


For the bipolar

$$R_{in,base} = (\beta + 1)(R_E + r_e)$$

$$R_{in,emitter} = \left(r_e + \frac{R_B}{\beta + 1} \right) \left(\frac{r_o + R_C}{r_o} \right) \cong \left(r_e + \frac{R_B}{\beta + 1} \right)$$

$$R_{in,collector} = r_o \left\{ 1 + g_m R_E \left(\frac{r_{\pi}}{r_{\pi} + R_E + R_B} \right) \right\} \cong r_o \{ 1 + g_m R_E \}$$



For the fet: THE EQUATIONS ARE THE SAME IF β IS INFINITE!

$$R_{in,gate} = open$$

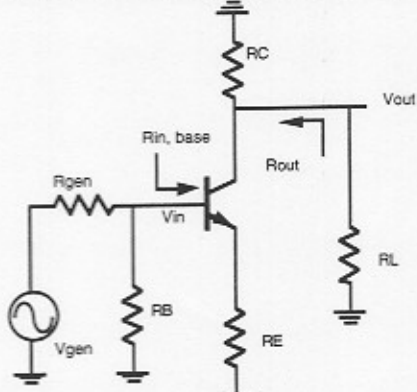
$$R_{in,source} = \left(\frac{1}{g_m} \right) \left(\frac{R_D + r_{ds}}{r_{ds}} \right) \cong \left(\frac{1}{g_m} \right)$$

$$R_{in,drain} = r_{ds} \{ 1 + g_m R_S \}$$

Be warned that in the equations above R_C , R_E and R_B are the equivalent resistances seen by the transistor. So in, the amplifier circuits below, THINK: "what is the effective resistance seen by the transistor?", before plugging into the equations.

Amplifier Stages

Common Emitter and Common Source



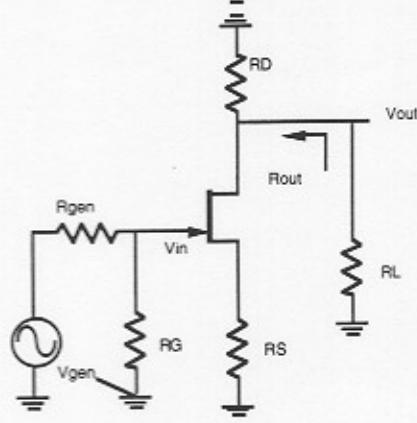
Common emitter:

$$R_{in} = R_B \parallel R_{in,base} = R_B \parallel ((\beta + 1)(R_E + r_e))$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{-(R_{Leq})}{r_e + R_E} = \frac{-(R_C \parallel R_L \parallel R_{out,collector})}{r_e + R_E}$$

$$R_{out} = (R_C \parallel R_{out,collector})$$



Common source:

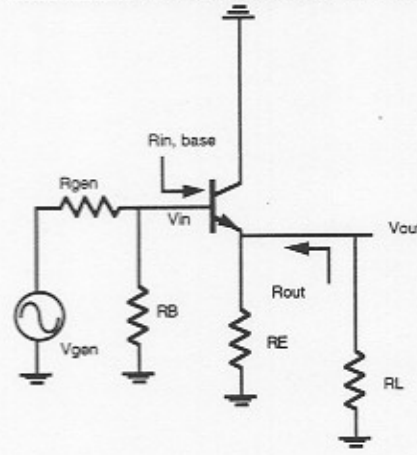
$$R_{in} = R_G$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{-(R_{Leq})}{1/g_m + R_S} = \frac{-(R_D \parallel R_L \parallel R_{out,drain})}{1/g_m + R_S}$$

$$R_{out} = R_D \parallel R_{out,drain}$$

Common Collector (emitter follower) and common drain (source follower)



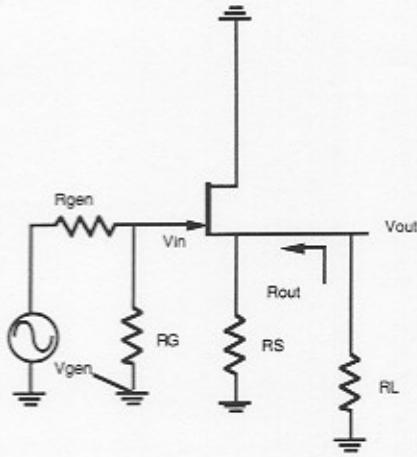
for the emitter follower:

$$R_{in} = R_B \parallel R_{in,base} = R_B \parallel ((\beta + 1)(R_E \parallel R_L + r_e))$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{R_{Leq}}{r_e + R_{Leq}} = \frac{R_E \parallel R_L}{r_e + R_E \parallel R_L}$$

$$R_{out} = R_E \parallel R_{in,emitter} = R_E \parallel \left(r_e + \frac{R_B}{\beta + 1} \right)$$



For the source follower:

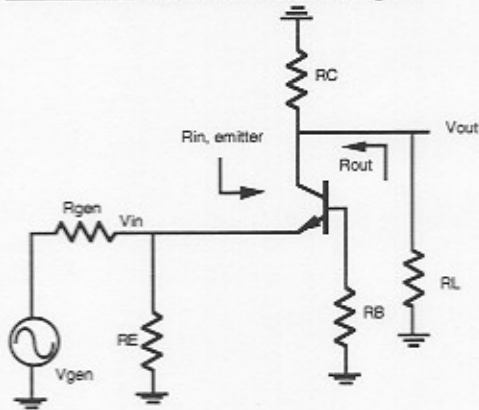
$$R_{in} = R_G$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{R_{Leq}}{1/g_m + R_{Leq}} = \frac{R_S \parallel R_L}{1/g_m + R_S \parallel R_L}$$

$$R_{out} = R_S \parallel R_{in,source} = R_S \parallel \frac{1}{g_m}$$

Common Base and common gate



For the common base:

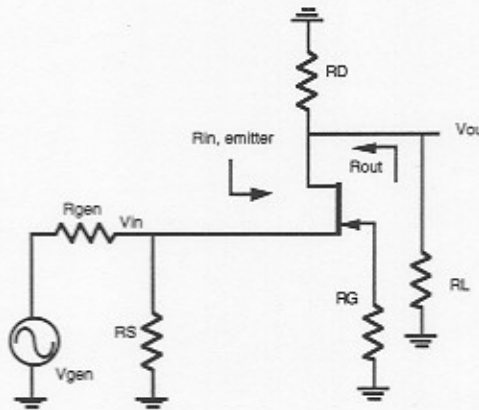
$$R_{in} = R_E \parallel R_{in,emitter}$$

$$R_{in,emitter} = \left(r_e + \frac{R_B}{\beta + 1} \right) \left(\frac{r_o + R_C \parallel R_L}{r_o} \right) \cong \left(r_e + \frac{R_B}{\beta + 1} \right)$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{R_{L,eq}}{R_{in,emitter}} = \frac{R_C \parallel R_L}{R_{in,emitter}}$$

$$R_{out} = (R_C \parallel R_{out,collector})$$



For the common gate

$$R_{in} = R_S \parallel R_{in,source}$$

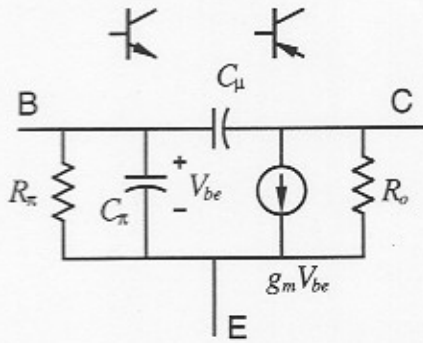
$$R_{in,source} = \left(\frac{1}{g_m} \right) \left(\frac{R_D \parallel R_L + r_{ds}}{r_{ds}} \right) \cong \left(\frac{1}{g_m} \right)$$

$$\frac{V_{in}}{V_{gen}} = \frac{R_{in}}{R_{in} + R_{gen}}$$

$$A_V = \frac{V_{out}}{V_{in}} = \frac{R_{L,eq}}{R_{in,source}} = \frac{R_D \parallel R_L}{R_{in,source}}$$

$$R_{out} = R_D \parallel R_{out,drain}$$

Basics: Transistor small-signal high-frequency models



Hybrid-pi model:

$$C_\mu = C_{bc}$$

$$C_\pi = C_{be} = C_{\pi,depl} + C_{\pi,diff}$$

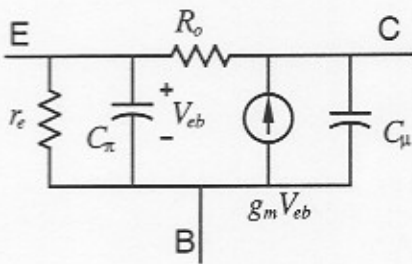
$$C_{\pi,diff} = g_m \tau_f$$

$$f_\tau = g_m / (2\pi(C_\pi + C_\mu))$$

$$g_m = \alpha / r_e = I_c / V_T$$

$$R_\pi = (\beta + 1)r_e$$

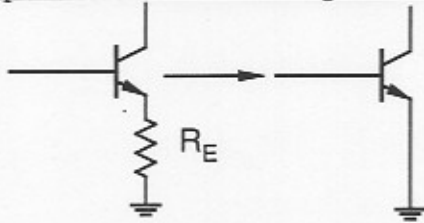
$$R_o = V_A / I_c$$



"T"-model

Makes common-base analysis much easier. Both these models are only approximate, being "good" up to f_t . Additionally R_{cb} is often an important parasitic (not discussed much in 137B).

Simplification of Emitter Degeneration



approximate only; check notes for bounds on validity

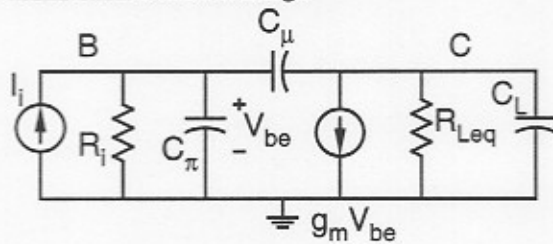
$$\tilde{C}_\pi = C_\pi (r_e / (r_e + R_E))$$

$$\tilde{g}_m = g_m (r_e / (r_e + R_E)) = \alpha / (r_e + R_E)$$

$$\tilde{C}_\mu = C_\mu$$

$$\tilde{R}_\pi = R_\pi ((r_e + R_E) / r_e) = (\beta + 1)(r_e + R_E)$$

Common-Emitter Stage



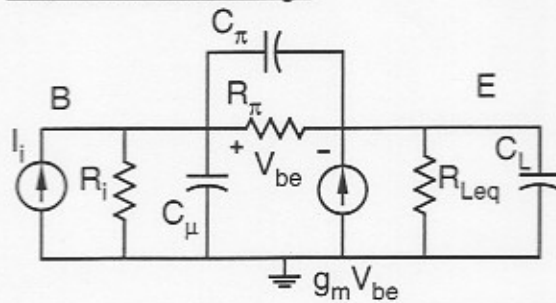
$$V_{out} / V_{gen} = (V_{out} / V_{gen})_{MB} \frac{1 + s\tau_{zero}}{1 + a_1 s + a_2 s^2}$$

$$a_1 = R_i (C_\pi + C_\mu (1 + g_m R_{Leq})) + R_{Leq} (C_\mu + C_L)$$

$$a_2 = R_i R_{Leq} (C_\mu C_L + C_\mu C_\pi + C_\pi C_L)$$

$$\tau_{zero} = -C_\mu / g_m$$

Emitter-Follower Stage



$$V_{out}/V_{gen} = (V_{out}/V_{gen})_{MB} \frac{1 + s\tau_{zero}}{1 + a_1s + a_2s^2}$$

given that $A_{vmb} = (r_e / (r_e + R_{Leq}))$:

$$a_1 = C_\pi (R_\pi \parallel (r_e \parallel R_{Leq} + R_i(1 - A_{vmb}))) + C_\mu (R_i \parallel \text{transistor input resistance}) + C_L (R_{Leq} \parallel \text{transistor output resistance})$$

$$a_2 = (R_i \parallel \text{transistor input resistance}) (R_{Leq} \parallel r_e) \times (C_\mu C_\pi + C_\mu C_L + C_L C_\pi)$$

$$\tau_{zero} = g_m / C_\pi$$

General Solutions of Problems: Nodal analysis (Know how to do this!)

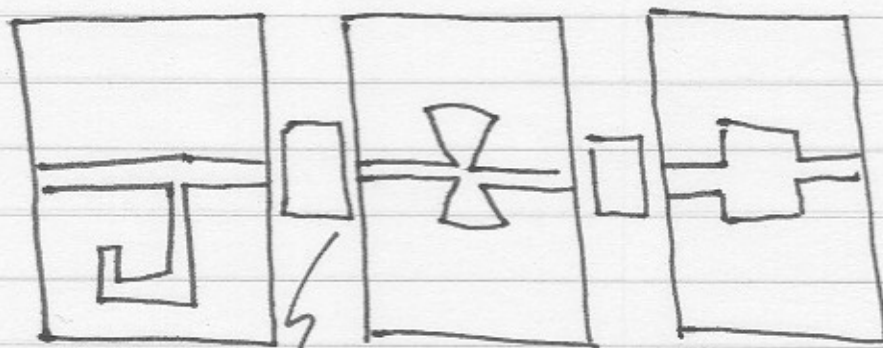
<p>1) Write the nodal equations (sum of the currents=0) at each circuit node, and put the resulting equations in matrix form (the Y's being various combinations of g_m's, $1/R$'s, and sC's):</p> $\begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & Y_{14} \\ Y_{21} & Y_{22} & Y_{23} & Y_{24} \\ Y_{31} & Y_{32} & Y_{33} & Y_{34} \\ Y_{41} & Y_{42} & Y_{43} & Y_{44} \end{bmatrix} \begin{bmatrix} V_1 = V_{in} \\ V_2 \\ V_3 \\ V_4 = V_{out} \end{bmatrix} = \begin{bmatrix} I_{in} \\ 0 \\ 0 \\ 0 \end{bmatrix}$	<p>2) Use Cramer's rule to solve:</p> $\begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & I_{in} \\ Y_{21} & Y_{22} & Y_{23} & 0 \\ Y_{31} & Y_{32} & Y_{33} & 0 \\ Y_{41} & Y_{42} & Y_{43} & 0 \end{bmatrix} = V_{out} \begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & Y_{14} \\ Y_{21} & Y_{22} & Y_{23} & Y_{24} \\ Y_{31} & Y_{32} & Y_{33} & Y_{34} \\ Y_{41} & Y_{42} & Y_{43} & Y_{44} \end{bmatrix}$
<p>3) This comes out as:</p> $\frac{V_{out}}{I_{in}} = s^m \frac{c_0 + c_1s + c_2s^2 + \dots}{d_0 + d_1s + d_2s^2 + \dots}$ <p>, which is divided through to get:</p> <p>(if present, m is the number of zeros, minus the number of poles, in the transfer function)</p>	<p>4)</p> $\frac{V_{out}}{I_{in}} = \left(\frac{V_{out}}{I_{in}} \right)_{mb} s^m \frac{1 + b_1s + b_2s^2 + \dots}{1 + a_1s + a_2s^2 + \dots}$ <p>...and the poles and zeroes are found by factoring the numerator and denominator. The separated-pole approximation, if applicable, makes this factoring easy.</p>
<p>5) To find the <i>impulse response</i>, do a partial-fraction expansion and then take the inverse LaPlace transform.</p>	<p>6) To find the sinusoidal frequency response, set $s = j\omega$.</p>

General Solutions of Problems: Method of Time Constants

	$\frac{V_{out}}{V_{gen}} = \left(\frac{V_{out}}{V_{gen}} \right)_{mb} \frac{1 + b_1 s + b_2 s^2 + \dots}{1 + a_1 s + a_2 s^2 + \dots}$ $a_1 = R_{11}^0 C_1 + R_{22}^0 C_2 + R_{33}^0 C_3 + R_{44}^0 C_4$ $a_2 = R_{11}^0 R_{22}^1 C_1 C_2 + R_{11}^0 R_{33}^1 C_1 C_3 + R_{11}^0 R_{44}^1 C_1 C_4 + R_{22}^0 R_{33}^2 C_2 C_3 + R_{22}^0 R_{44}^2 C_2 C_4 + R_{33}^0 R_{44}^3 C_3 C_4$ <p>notethat $R_{xx}^0 R_{yy}^x = R_{xx}^y R_{yy}^0$</p>
	$R_{xx}^0 = R_x \parallel (r_e \parallel R_{Leq} + R_i (1 - A_{vmb}))$ $A_{vmb} = (R_{Leq} / (r_e + R_{Leq}))$
	$R_i = R_x \parallel R_x$ $R_{yy}^0 = R_i (1 + g_m R_{Leq}) + R_i$

This brings me to my second major point:

classical microwave design looks like this.



transistor die
microwave matching networks on alumina.

Classical Microwave Smith-Chart design works very well for this.

But industry is moving increasingly to IC's

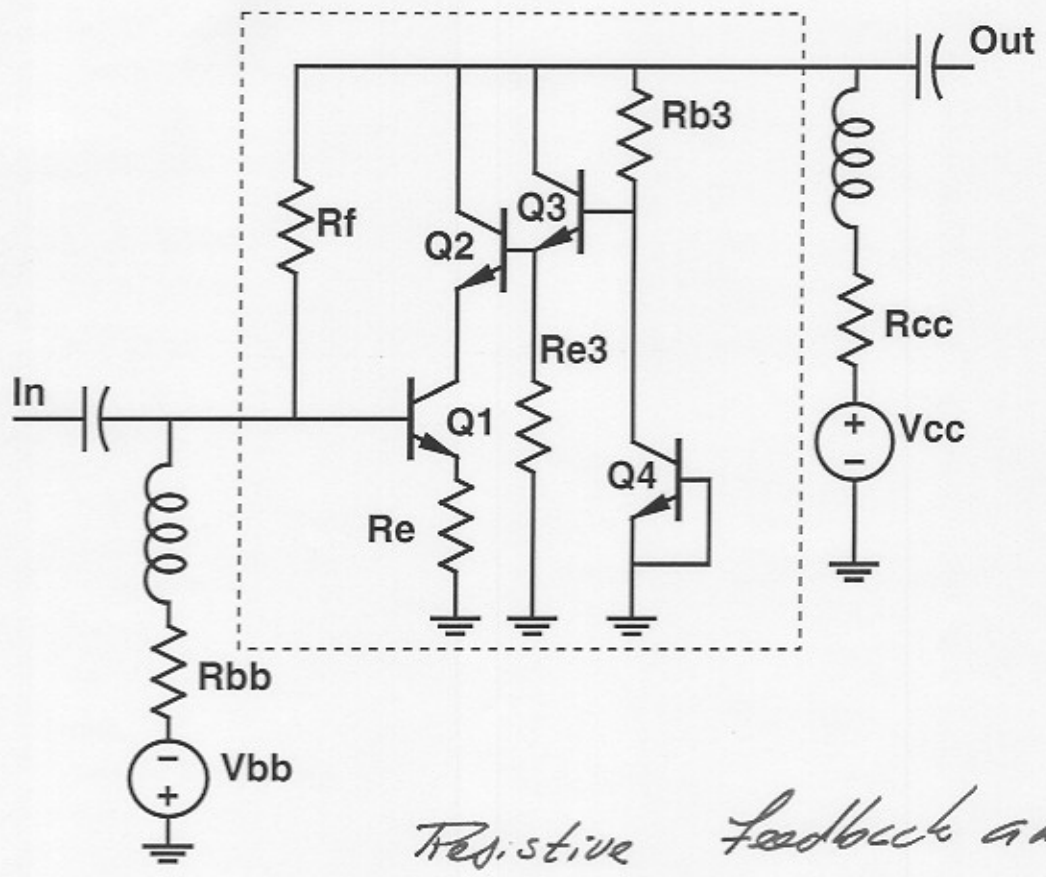
- Some look just like the prior example.

- But many look like the examples attached on the next 2 pages.

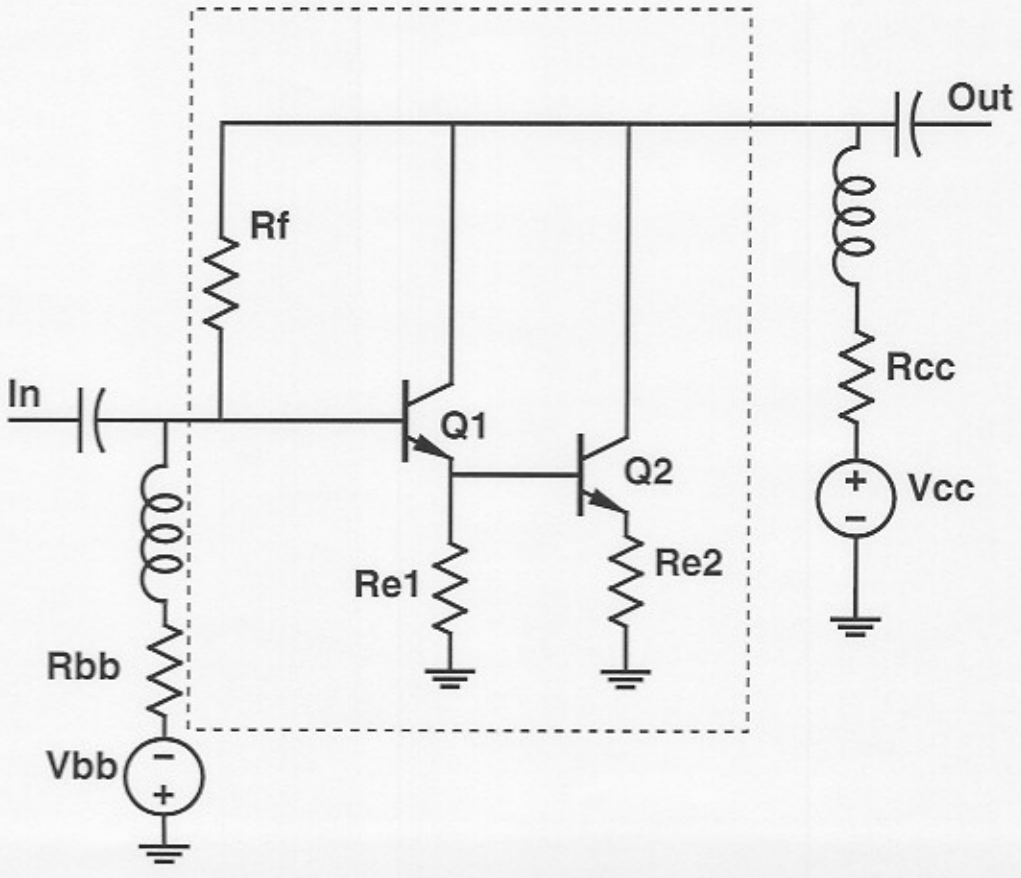
For these, we will need a mixture of classical microwave design and analog IC design methods. Both are important.

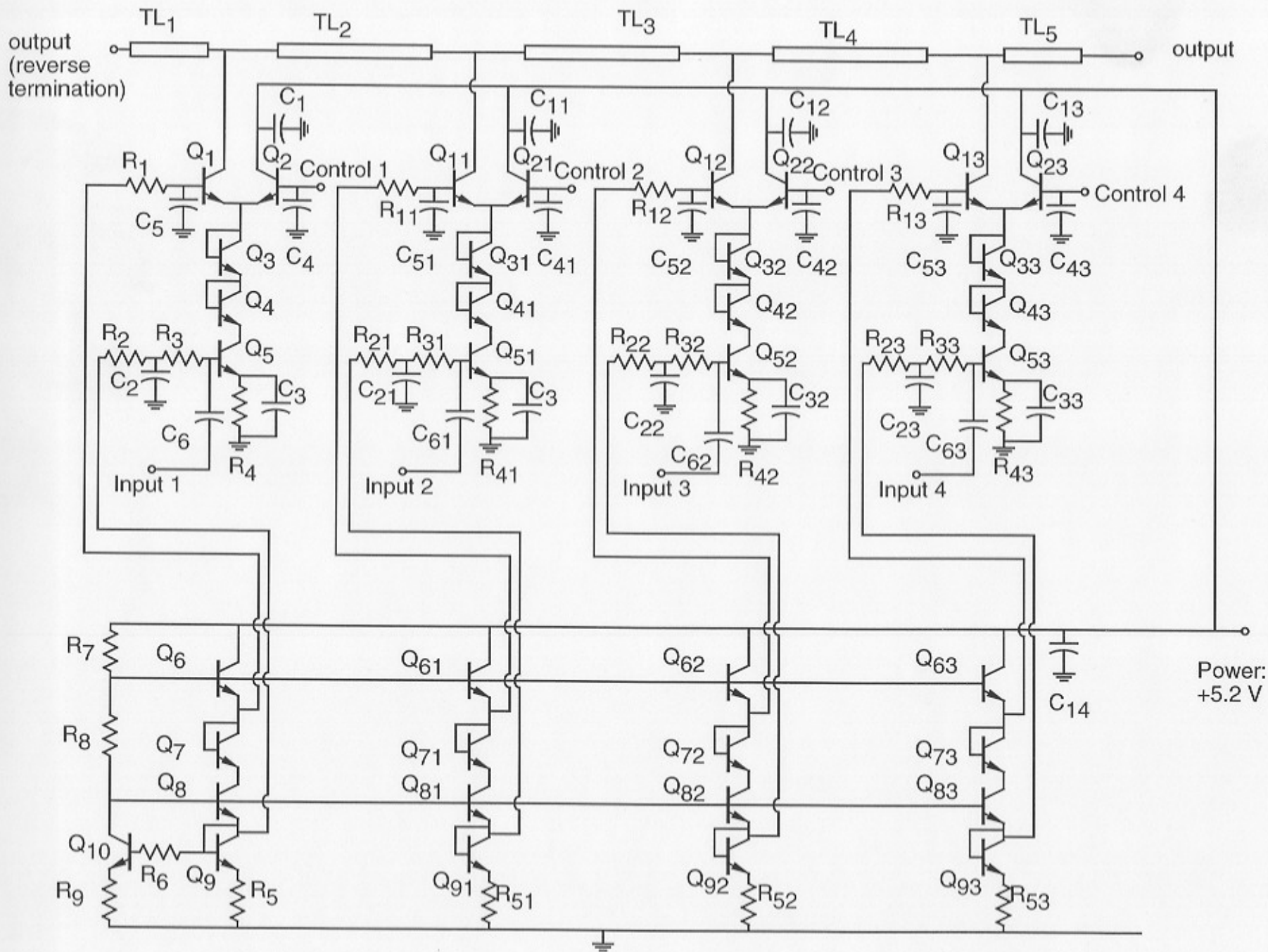
An excellent reference for "analog IC design" is

Grey and Meyer "Analog Integrated Circuits"



Resistive feedback amplifiers...





Manchester Active Switch.

Many of these circuits can and do operate over bandwidths as high as 30 GHz.

→ "Analog Circuit design" techniques are perfectly applicable at microwave frequencies!

Methods for analog design:

* Nodal Analysis

* Method of Time Constants.

Let's review these below.

Please refer extensively to the

cr. 6 sheets on pp. 2-7.

Nodal Analysis

General Solutions of Problems: Nodal analysis (Know how to do this!)

<p>1) Write the nodal equations (sum of the currents=0) at each circuit node, and put the resulting equations in matrix form (the Y's being various combinations of g_m's, $1/R$'s, and sC's):</p> $\begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & Y_{14} \\ Y_{21} & Y_{22} & Y_{23} & Y_{24} \\ Y_{31} & Y_{32} & Y_{33} & Y_{34} \\ Y_{41} & Y_{42} & Y_{43} & Y_{44} \end{bmatrix} \begin{bmatrix} V_1 = V_{in} \\ V_2 \\ V_3 \\ V_4 = V_{out} \end{bmatrix} = \begin{bmatrix} I_{in} \\ 0 \\ 0 \\ 0 \end{bmatrix}$	<p>2) Use Cramer's rule to solve:</p> $\begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & I_{in} \\ Y_{21} & Y_{22} & Y_{23} & 0 \\ Y_{31} & Y_{32} & Y_{33} & 0 \\ Y_{41} & Y_{42} & Y_{43} & 0 \end{bmatrix} = V_{out}$ $\begin{bmatrix} Y_{11} & Y_{12} & Y_{13} & Y_{14} \\ Y_{21} & Y_{22} & Y_{23} & Y_{24} \\ Y_{31} & Y_{32} & Y_{33} & Y_{34} \\ Y_{41} & Y_{42} & Y_{43} & Y_{44} \end{bmatrix}$
<p>3) This comes out as:</p> $\frac{V_{out}}{I_{in}} = k s^m \frac{c_0 + c_1 s + c_2 s^2 + \dots}{d_0 + d_1 s + d_2 s^2 + \dots}$ <p>, which is divided through to get:</p> <p>(if present, m is the number of zeros, minus the number of poles, in the transfer function)</p>	<p>4)</p> $\frac{V_{out}}{I_{in}} = \left(\frac{V_{out}}{I_{in}} \right)_{mb} s^m \frac{1 + b_1 s + b_2 s^2 + \dots}{1 + a_1 s + a_2 s^2 + \dots}$ <p>...and the poles and zeroes are found by factoring the numerator and denominator. The separated-pole approximation, if applicable, makes this factoring easy.</p>
<p>5) To find the <i>impulse response</i>, do a partial-fraction expansion and then take the inverse LaPlace transform.</p>	<p>6) To find the sinusoidal frequency response, set $s = j\omega$.</p>

lets give an example of this...

Before we do, what is

"the separated-pole approximation"?

$$\text{If we have } H(s) = \frac{N(s)}{1 + a_1 s + a_2 s^2}$$

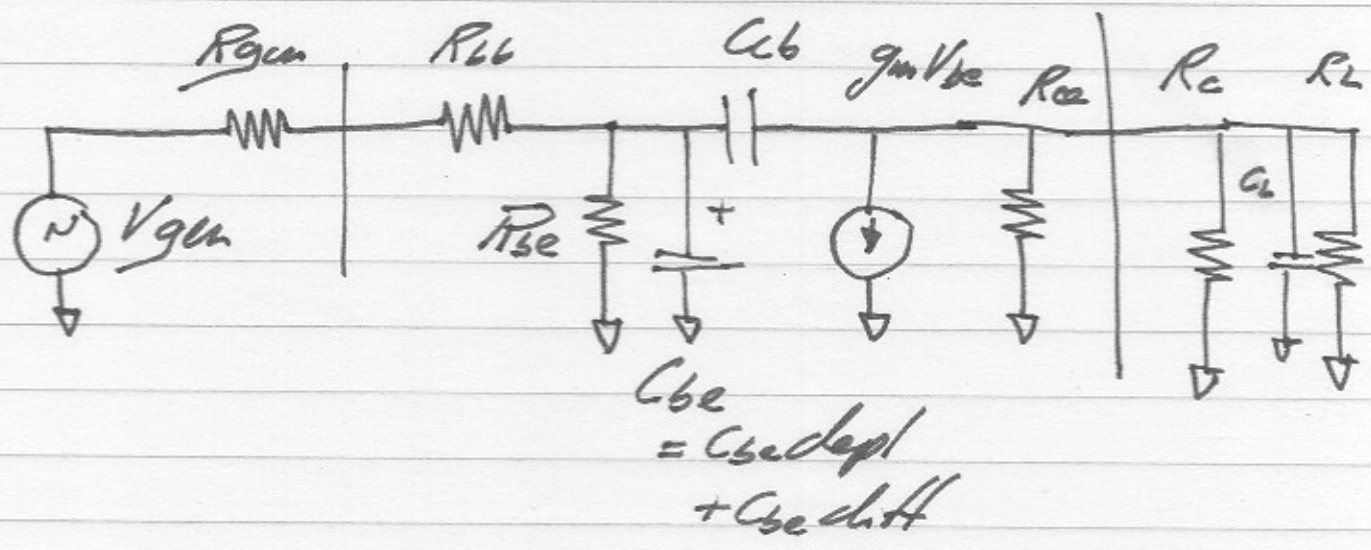
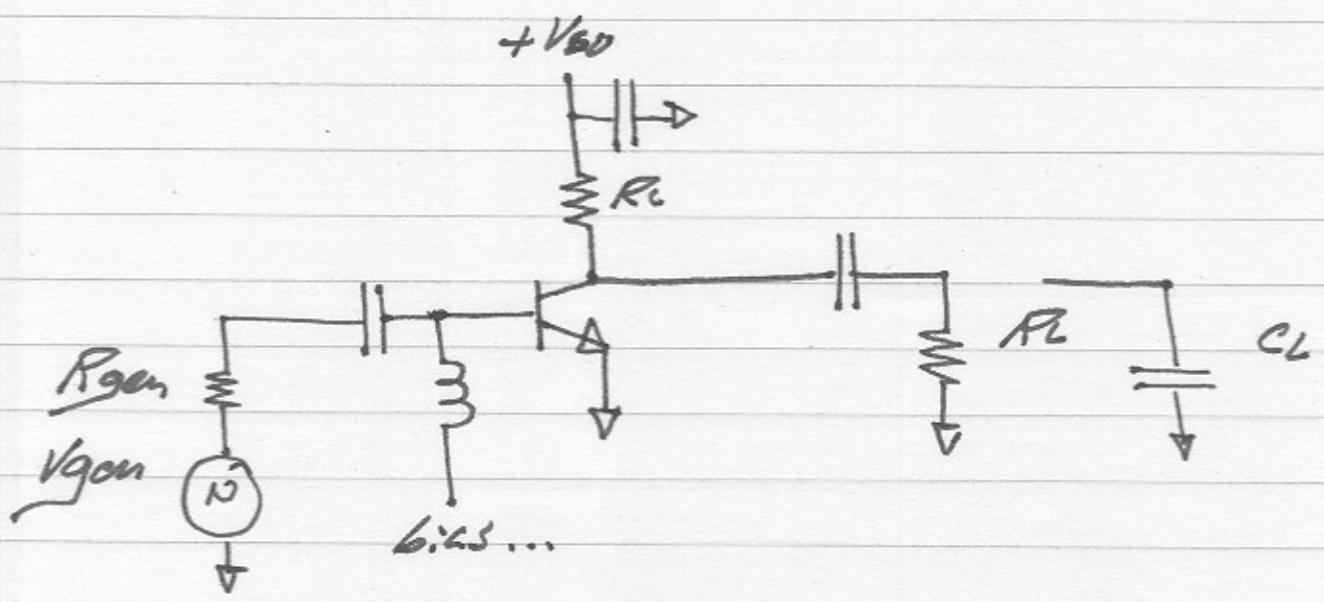
$$H(s) \cong N(s) \cdot \frac{1}{1 + a_1 s} \cdot \frac{1}{1 + a_2 s/a_1}$$

$$= N(s) \frac{1}{1 + (a_1 + \frac{a_2}{a_1})s + a_2 s^2}$$

If $a_1 \gg a_2/a_1$

This a valuable method of finding the roots (poles) of a transfer function without resorting to the quadratic equation...

Example: Common-Emitter stage:



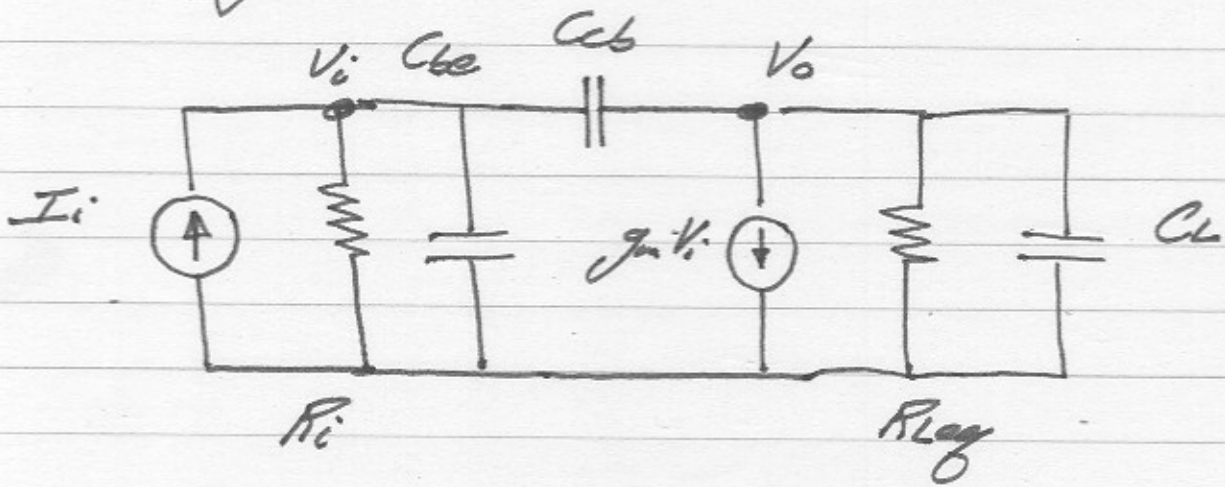
note the simplicity of the models assumed.

I will convert V_{gen} , R_{bb} & R_{gen} with a Thevenin-Norton Transformation.

and write $R_i = R_{be} \parallel (R_{bb} + R_{gen})$

$$R_{eq} = R_{ce} \parallel R_L \parallel R_C$$

\Downarrow $I_i = V_{gen} / (R_{bb} + R_{gen})$



So now we write $\sum I = 0$ at each node where the voltage is unknown (V_i & V_o)

$$\begin{cases} V_i \left(\frac{1}{R_i} + \beta C_{be} + \beta C_{cb} \right) + V_o (-\beta C_{cb}) = I_i \\ V_i (-\beta C_{cb} + g_m) + V_o \left(\frac{1}{R_{eq}} + \beta C_{cb} + \beta C_L \right) = 0 \end{cases}$$

we then in general write this in matrix form to solve...

$$\begin{bmatrix} I_i \\ 0 \end{bmatrix} = \begin{bmatrix} \left(\frac{1}{R_i} + \beta C_{be} + \beta C_{cb} \right) & -\beta C_{cb} \\ (-\beta C_{cb} + g_m) & \left(\frac{1}{R_{eq}} + \beta C_{cb} + \beta C_L \right) \end{bmatrix} \begin{bmatrix} V_i \\ V_o \end{bmatrix}$$

and use Cramer's rule to solve:

$$\rightarrow \frac{V_{out}}{V_{gen}} = \frac{V_{out}}{V_{gen}} \bigg|_{MS} \frac{1 + s T_{zero}}{1 + a_1 s + a_2 s^2}$$

$$a_1 = R_i (C_{be} + C_{cb} (1 + g_m R_{leg})) + R_{leg} (C_{cb} + C_L)$$

$$a_2 = R_i R_{leg} (C_{cb} C_{be} + C_{cb} C_L + C_{be} C_L)$$

$$T_{zero} = -C_{cb} / g_m$$

note that I have glossed over ~ 2 pages of algebra to obtain this.

We now come to use the separated pole approximation:

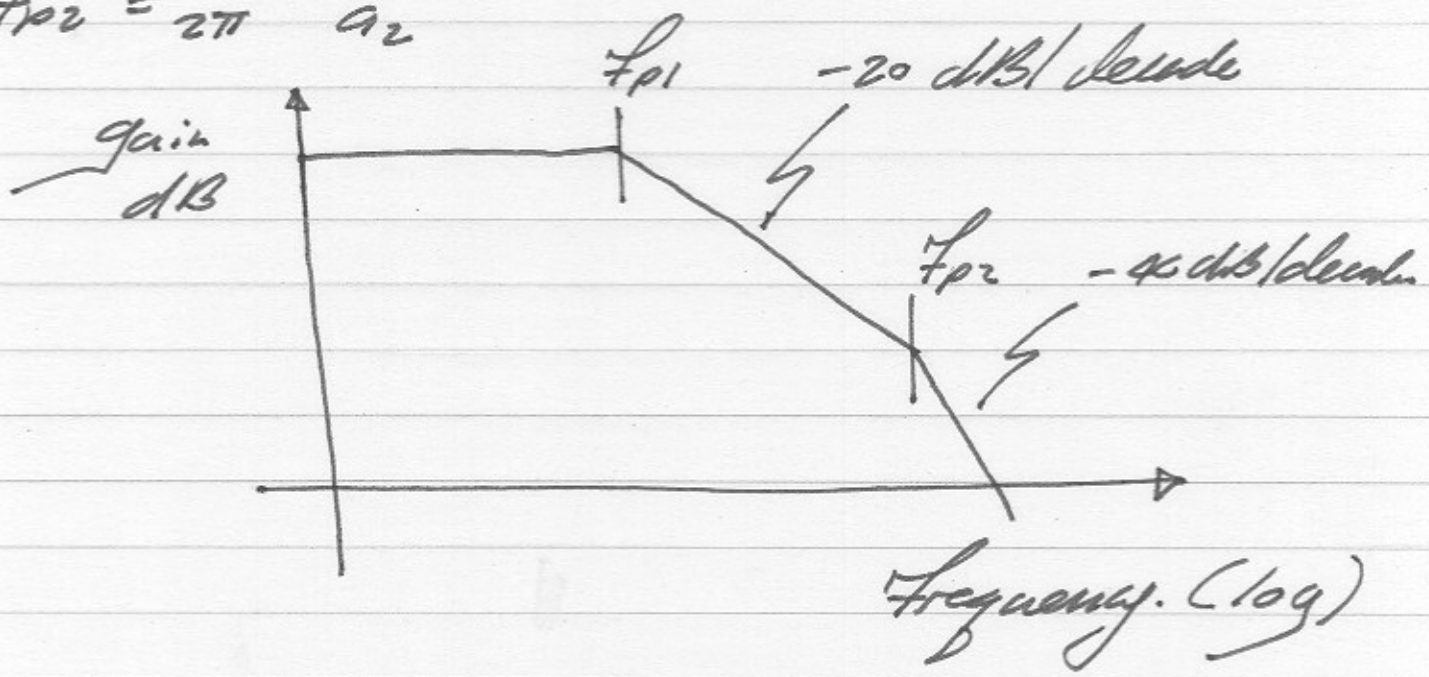
$$\frac{V_{out}}{V_{gen}} \approx \frac{V_{out}}{V_{gen}} \bigg|_{MS} \frac{1 + s T_{zero}}{(1 + a_1 s) (1 + \frac{a_2}{a_1} s)}$$

which gives pole frequencies of

$$f_{p1} = \frac{1}{2\pi a_1}$$

and

$$f_{p2} = \frac{1}{2\pi} \frac{a_1}{a_2}$$



Look at the answer

$$f_{p1} = \frac{1}{2\pi a_1}$$

$$a_1 = R_i \underbrace{(C_{be} + C_{cb}(1 + g_m R_{eq}))}_{\text{input RC time constant } T_{in}} + R_{eq} \underbrace{(C_{cb} + C_L)}_{\text{output RC time constant } T_{out}}$$

$$T_{in} = R_i C_{equiv}$$

where

$$C_{equiv} = C_{be} + C_{cb} \underbrace{(1 + g_m R_{eq})}_{1 - A_v}$$

this is

the familiar Miller Approximation.

OK so what have we accomplished:

- obtained gain-bandwidth relationships for an elementary c.f. stage.

- I hope this was just a reminder

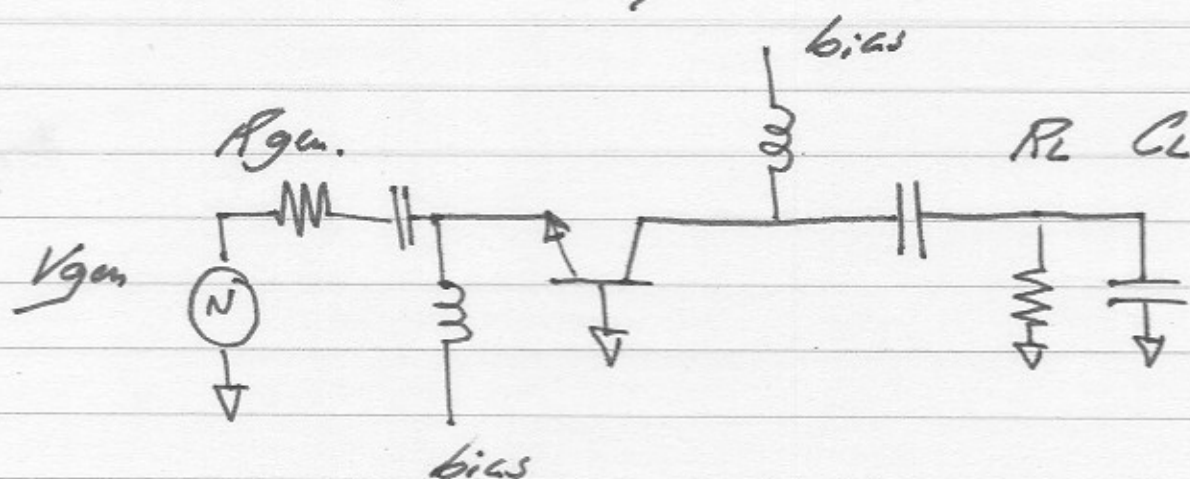
* Reviewed HOW to generally solve frequency response problems



This is the message. Although a simple procedure, one has to do this all the time in evaluating complex IC's.

Next ExampleSimplified Common-base stage

You are warned that here we ignore R_{bb} which is a very dangerous assumption.

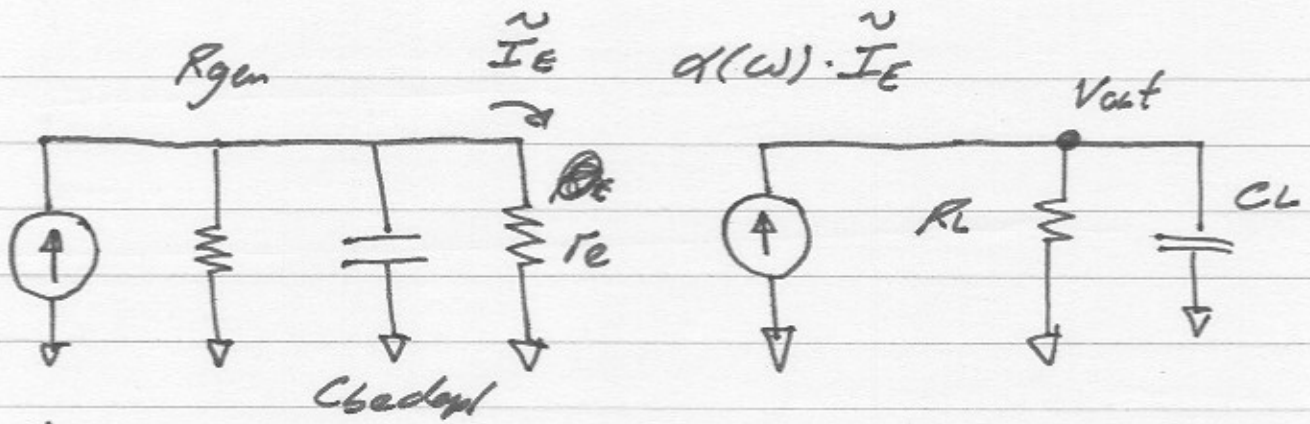


we will use the T-model

Simplify: $R_{\text{bb}} \rightarrow 0$ dangerous!

$R_{\text{ex}} \rightarrow 0$ (easy to fix)

$R_{\text{CE}} \rightarrow +\infty$ reasonable.



$I_i = V_{gen} / R_{gen}$

by inspection:

$$\frac{V_{out}}{V_{gen}} \approx \frac{V_{out}}{V_{gen} \cdot \frac{1}{1 + j\omega C_{bedep1} (r_e || R_{gen})}}$$

- $\frac{1}{1 + j\omega (\tau_b + \tau_c)}$

- $\frac{1}{1 + j\omega R_L C_L}$

... The third term is usually dominant again, one is warned of the danger of ignoring $R_L C_L$.

What about the emitter follower stage?

- We will find this hard: lets
learn the method of time
constants first!

What about Field-Effect Transistors?

The models of T_n and T_p
are basically similar - only the
notation changes ($C_{be} \rightarrow C_{gs}$ etc)